



DISSIPATIVE OPTICAL SOLITONS (DOSs): TOWARDS NANOSIZES AND ATTODURATIONS

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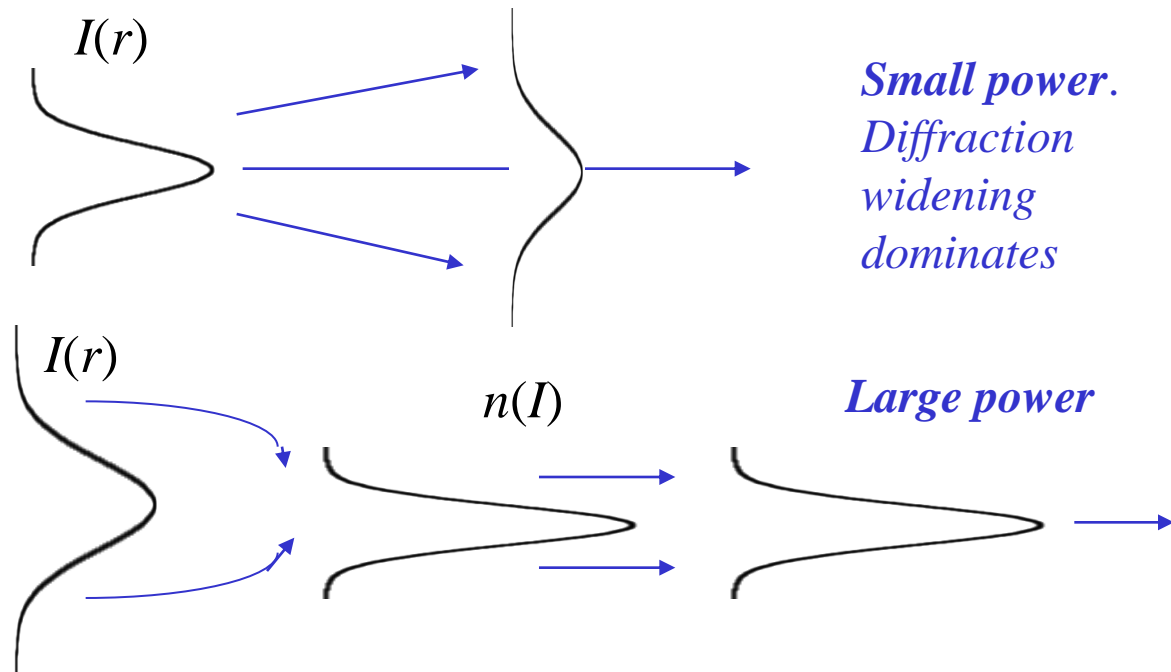
OUTLINE

DOS = Dissipative Optical Soliton, or Autosoliton

1. Introduction. Optical Soliton History, Terminology, Motivation
2. “Standard” envelope DOSs in wide-aperture driven nonlinear interferometers and lasers with saturable absorption
 - Pseudo-quantum features (discrete spectrum of main parameters)
 - Internal structure (topology of energy flows)
 - Weak and strong interaction of DOSs
 - Symmetry and Eulerian mechanics
3. Discrete dissipative solitons, including knot solitons.
4. Extreme DOSs
 - “Nanosolitons” – nanosized DOSs in driven molecular chains
 - “Attosolitons” – subfemtosecond DOSs on the basis of self-induced transparency
5. Conclusion

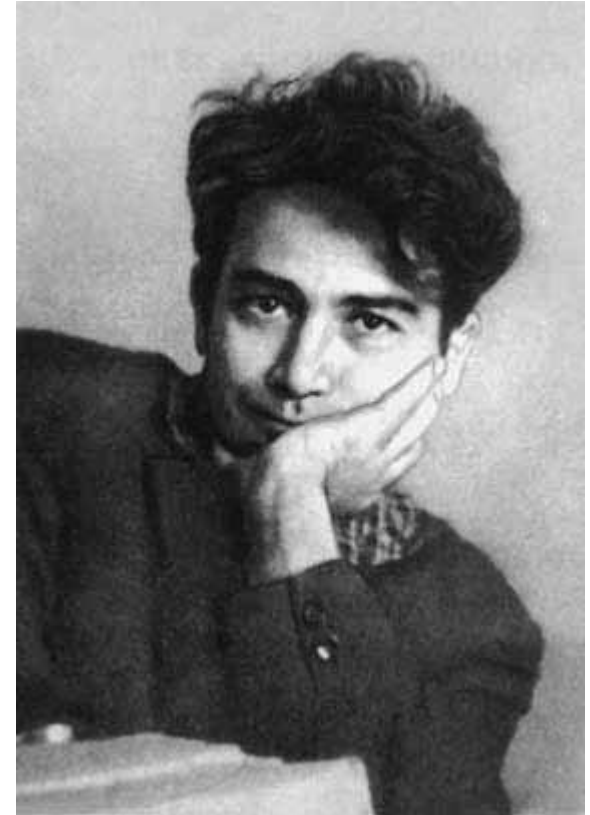
Conservative Spatial Optical Solitons (1962)

Self-focusing (including self-trapping) of optical beams in transparent media with intensity-dependent refractive index $n(I)$



Nonlinear focusing and their stabilization due to the balance with diffraction widening (below optical damage threshold)

Theory: Talanov, Litvak; Chiao, Garmire, Townes; Akhmanov, Sukhorukov, Khokhlov; Zakharov; ...

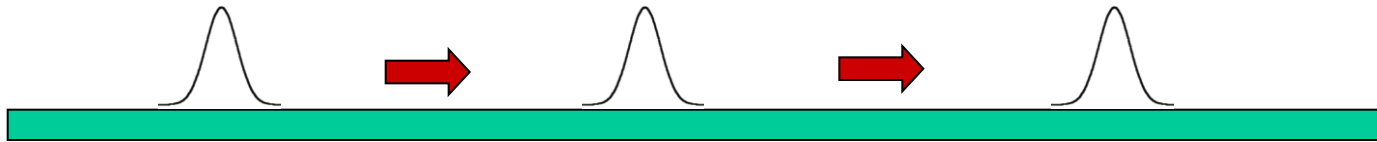


G.A. Askaryan
1928-1997

Are similar solitons possible with nonlinearity of gain and absorption?

Spatio-Temporal Analogy

Temporal solitons in a single-mode fiber [A. Hasegawa, F. Tappert, 1973]



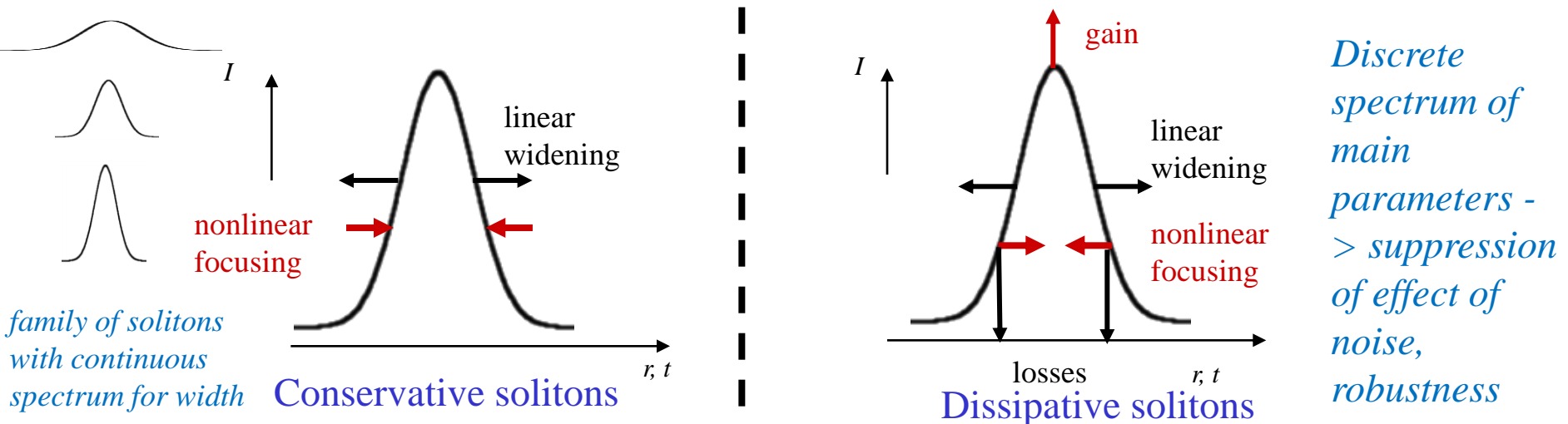
Optical pulse widening due to frequency dispersion is compensated by nonlinear focusing

Experiment: Mollenauer et al.

In Russia: Nonlinear fibre optics, A.M. Prokhrov, E.M. Dianov et al., ...

2002: Commercial solitonic communication line, DWDM, Australia, 3, 875 km, 1.6 Tbit/s

Yu.S. Kivshar, G.P. Agrawal. Optical Solitons. From Fibers to Photonic Crystals, Acad.Pr, 2003



AUTOSOLITON (*dissipative soliton*) is a stable field structure localized due to nonlinear balance of energy input and output in homogeneous or periodically modulated non-conservative medium or system.

Motivation for research of DOSs

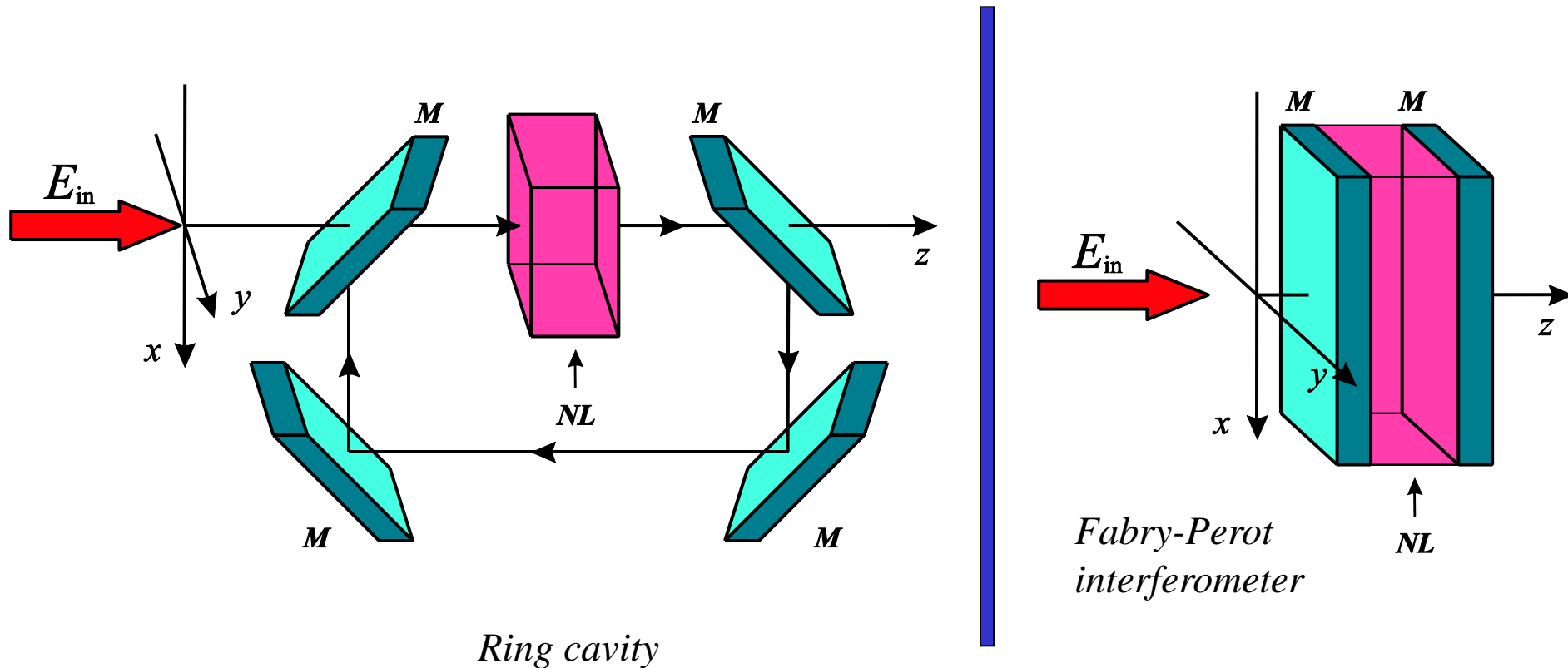
- For science: Wave-particle dualism.

If dissipative solitons corresponds to “particles”, what are features of these particles?

The diversity of the dissipative optical solitons and unusual type of their features.

- For applications: Coding information by solitons, and then to record, store, and process information. What are the features of dissipative solitons promising for these applications?
- Tendencies in research of dissipative optical solitons.

Driven Wide-Aperture Nonlinear Interferometer



Field phase and frequency are determined by holding radiation

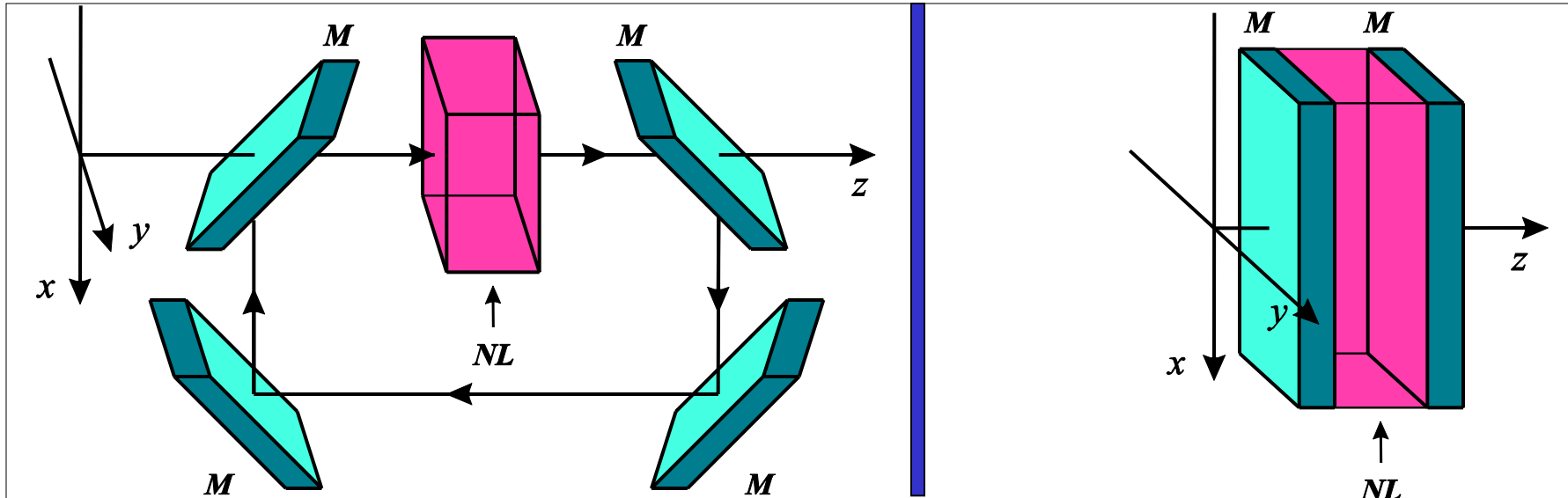
M – mirrors

NL – medium with optical nonlinearity

E_{in} – amplitude of coherent holding radiation

*[Rosanov, Semenov, 1980] –
modulational instability (filamentation);
[Rosanov, Semenov, Khodova,
1983, 1988, ...] – spatial DOSs*

Wide-Aperture Laser with Saturable Absorption



M – mirrors

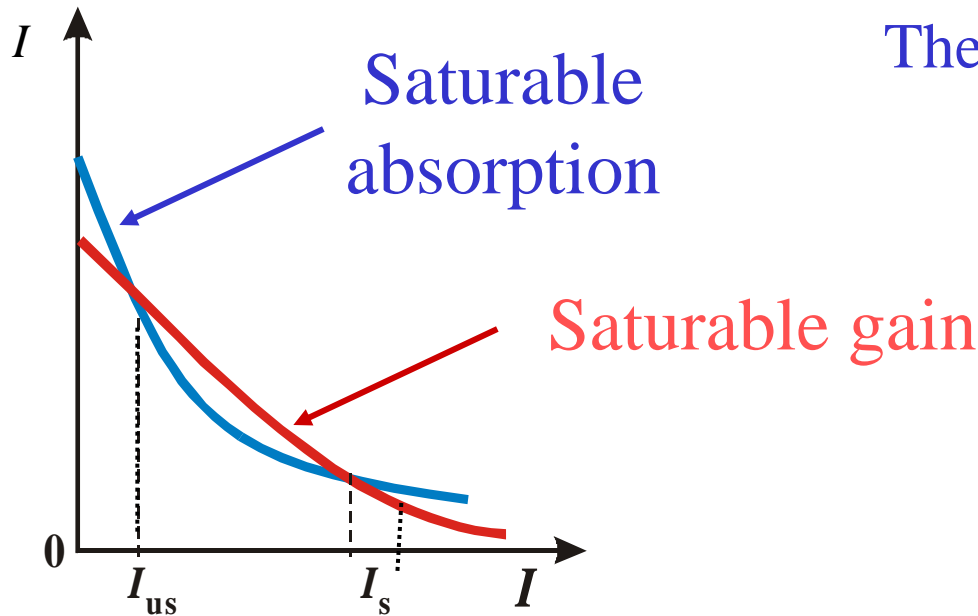
NL – medium with nonlinear gain and absorption

Envelope “laser solitons”:
[Rosanov, S. Fedorov, 1992]

No coherent holding radiation, incoherent pump (laser schemes)

Zero background. Free phase, frequency as eigenvalue.

Bistability and Laser Solitons

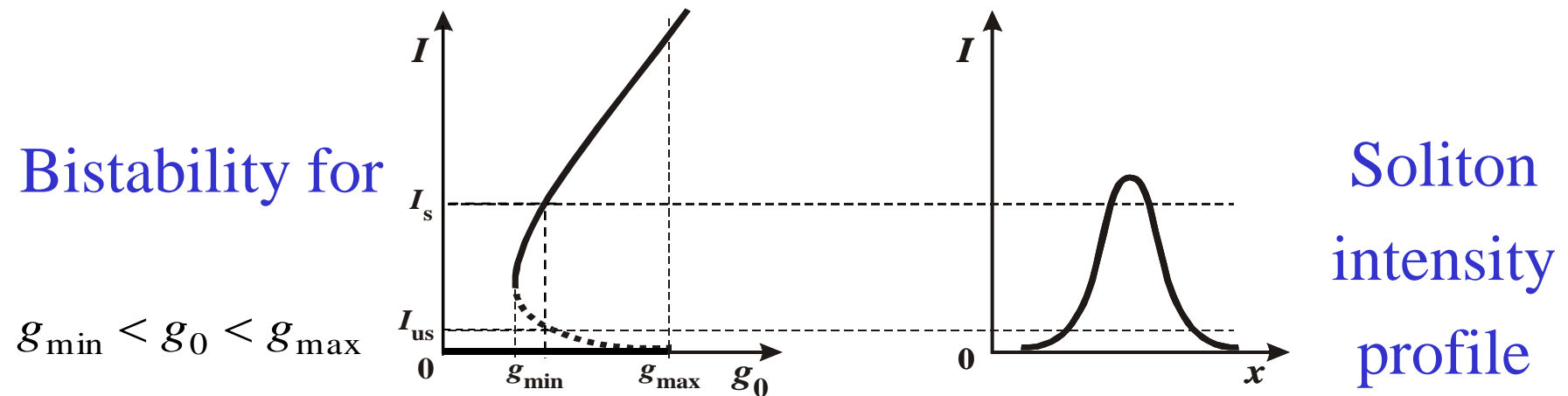


The case of instantaneous nonlinearity
(class A lasers)

$$\longleftarrow g_0 = \text{const}$$

g_0 - small-signal gain

$I = |E|^2$ - intensity



Hard excitation of lasing by external beam-pulse

Governing Equation for Laser Schemes

$$\frac{\partial E}{\partial \zeta} = (i + d) \Delta_D E + E f(|E|^2)$$

$$\tilde{E} \equiv \text{Re}[E \exp(ikz - i\omega t)]$$

$$I = |E|^2 - \text{intensity}$$

$$D = 1, 2, 3.$$

E – electric field envelope (mean-field approximation);

ζ – evolution variable (time or longitudinal co-ordinate);

d – effective diffusion coefficient (angular selectivity);

$$\Delta_D = \nabla_D^2 \quad - D\text{-dimensional Laplacian,} \quad \Delta_2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$f(|E|^2)$ – medium nonlinearity + linear losses

$$f(|E|^2) = -1 + \frac{g_0}{1 + |E|^2} - \frac{a_0}{1 + b|E|^2}$$

gain *absorption*

Class A laser with
fixed polarization

[Suchkov 1966]

Governing Equation for Interferometers

$$\frac{\partial E}{\partial t} = i\Delta_D E + E f(|E|^2) + E_{in}$$

E – electric field envelope (mean-field approximation);

E_{in} – holding radiation amplitude;

$\Delta_D = \nabla_D^2$ – D -dimensional Laplacian, $D = 1$ or 2 .

$f(|E|^2)$ – medium nonlinearity + linear losses + detuning

For interferometers with Kerr nonlinearity (2D)

[Lugiato, Lefever 1987]

$$\frac{\partial E}{\partial t} = \underset{\substack{\uparrow \\ \text{diffraction}}}{i\Delta_{\perp}} E + \underset{\substack{\uparrow \\ \text{refractive} \\ \text{index} \\ \text{nonlinearity}}}{i|E|^2} E - \underset{\substack{\uparrow \\ \text{linear} \\ \text{losses}}}{(1 + i\Theta)} E + \underset{\substack{\nwarrow \\ \text{holding} \\ \text{radiation}}}{E_{in}}$$

DOSs Quasi-Quantum Features

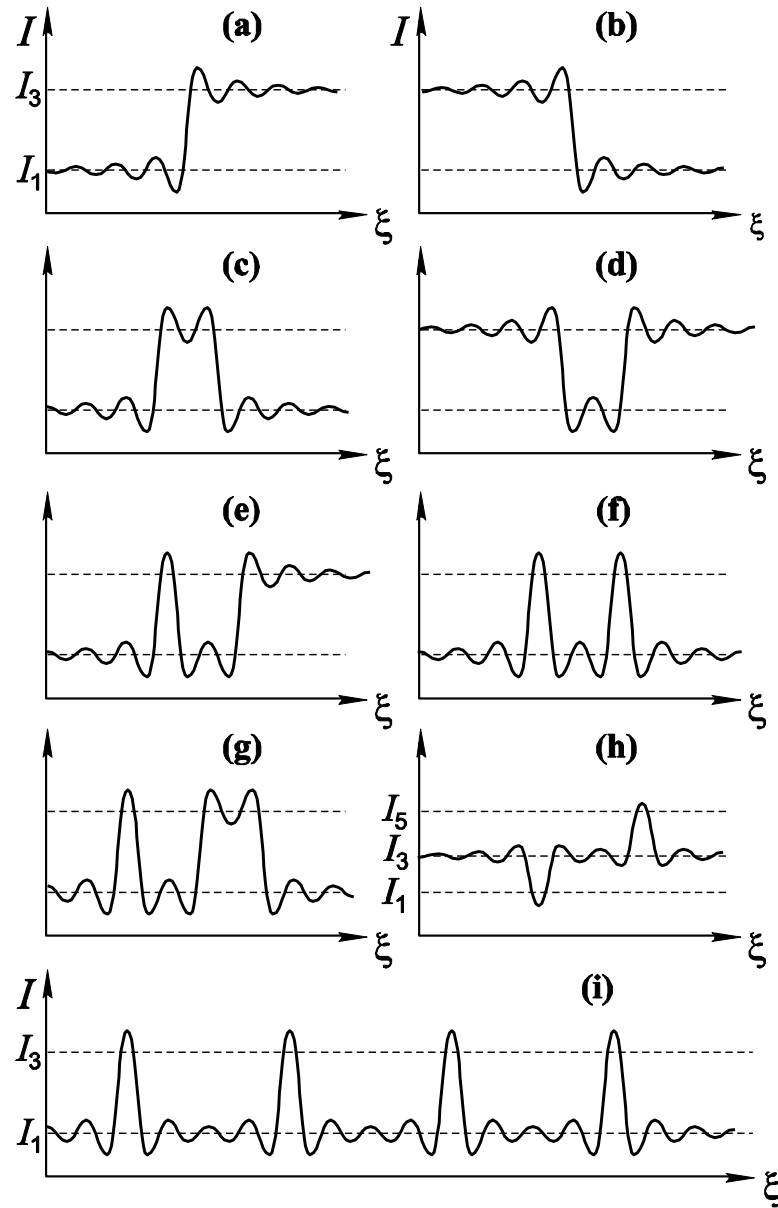
NONLINEAR INTERFEROMETER (1D)

Exactly solvable model: step-wise
nonlinearity [Rosanov, 1992]

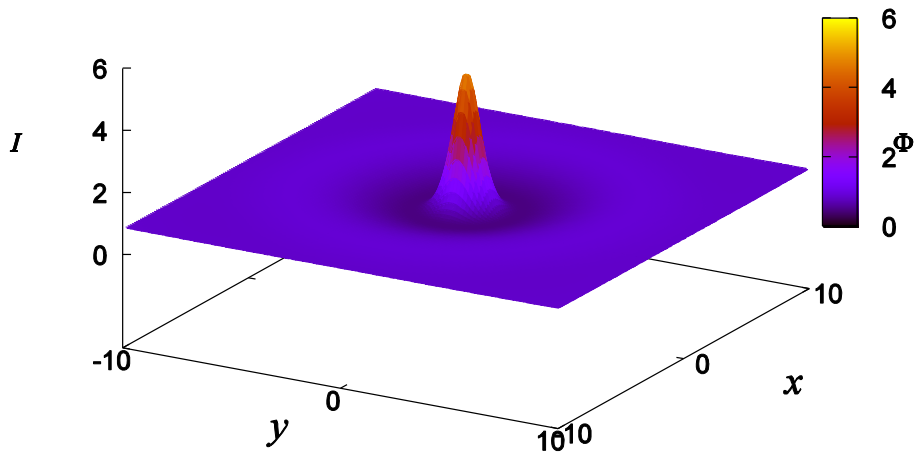
Discrete spectrum of single DOS width
– quasi-atom

Discrete spectrum of inter-soliton
distance for soliton pairs –
quasi-molecule

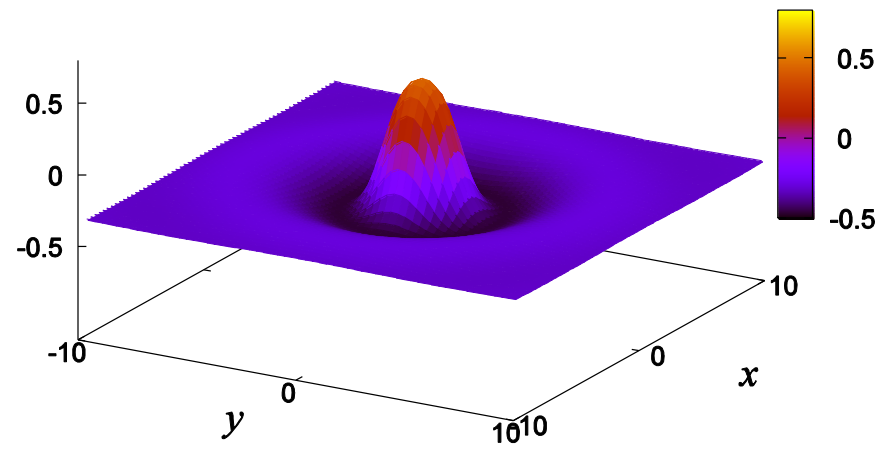
Band structure for periodic chains of
solitons – quasi-solid



Single 2D-Soliton in Interferometer



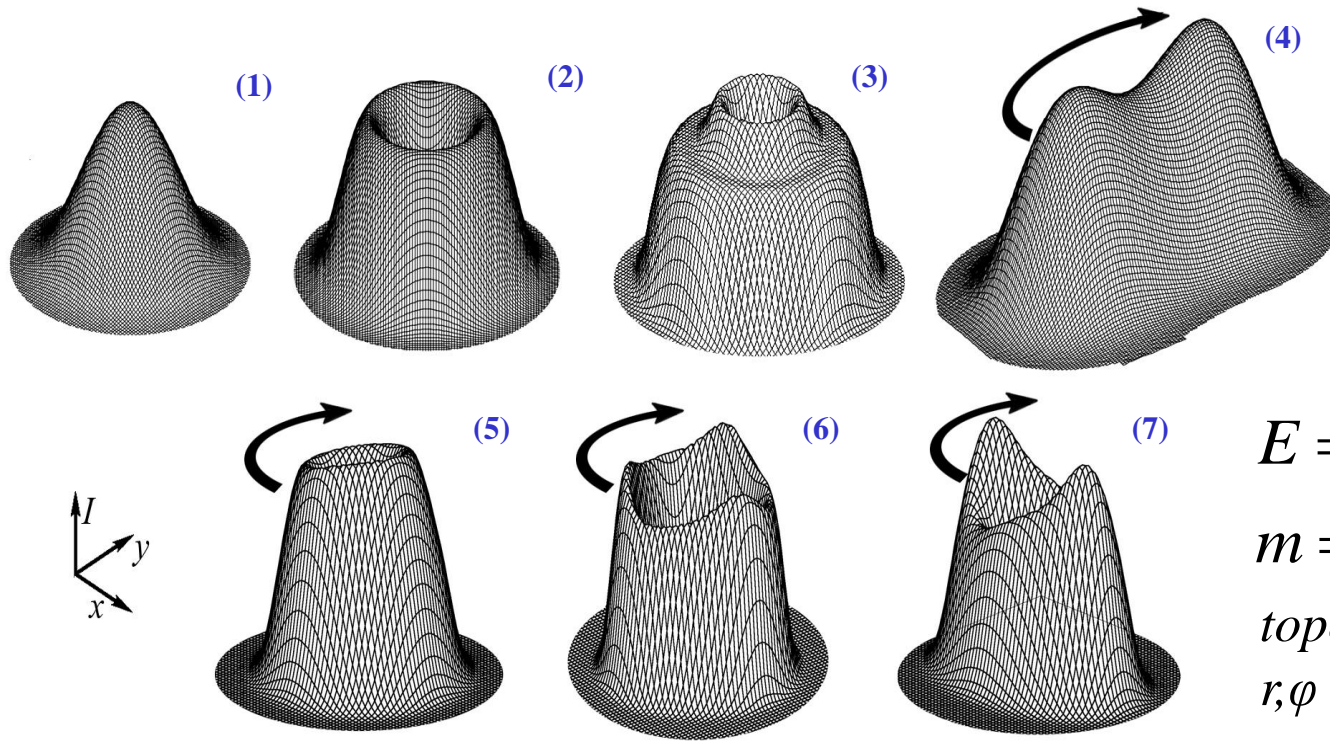
Intensity



Phase

2D-DOSs in lasers with saturable absorption

Transverse intensity distributions:



1 – 3:

$$E = A(r) \exp(im\varphi - i\alpha t)$$

$$m = 0, \pm 1, \pm 2 -$$

topological charge,

r, φ – polar coordinates,

α – frequency shift

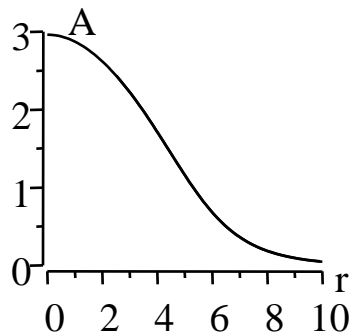
1 – fundamental DOS (regular wavefront); 2, 3, 5-7 – vortex DOSs;
4, 6, 7 – rotating DOSs; 5 – Moon-like motion

Soliton Internal Structure

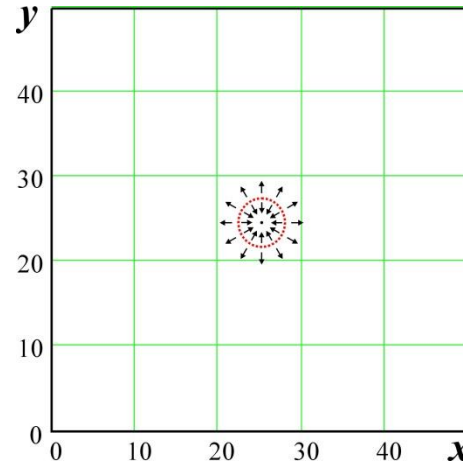
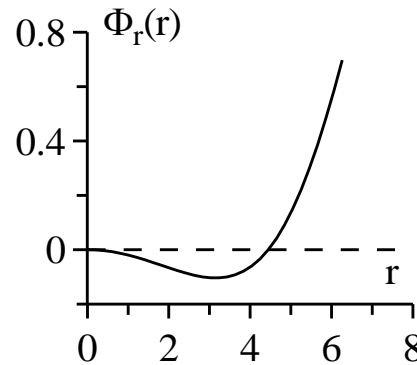
Single symmetric laser solitons: radial profiles and energy flows

Symmetry to rotation at any angle

$$A = |A(r)|$$



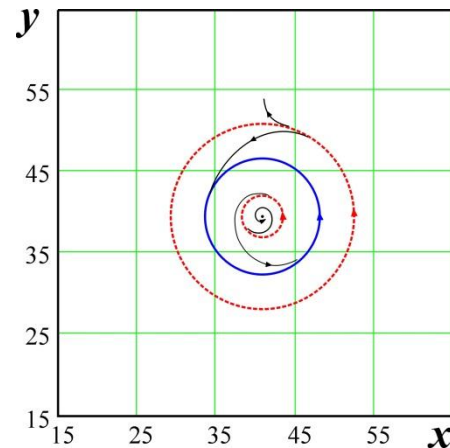
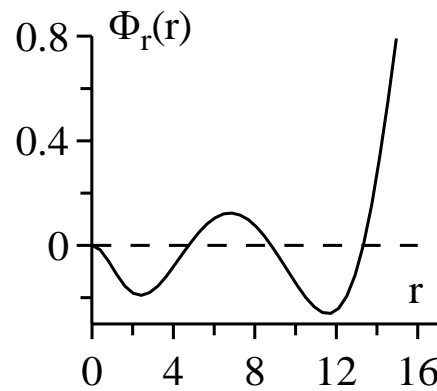
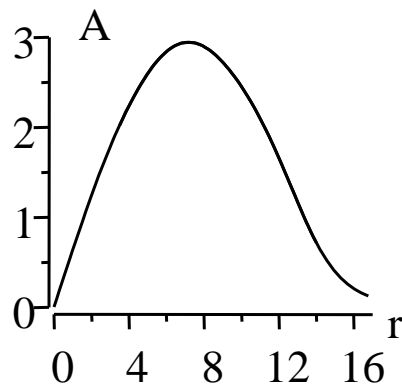
$$\Phi = \Phi_r(r) + m\varphi$$



Energy flows
portrait

$$m = 0$$

$$\mathbf{S}_{\perp} \sim I \nabla_{\perp} \Phi$$



$$m = 1$$

Focus,
one “stable” and
two “unstable”
limit cycles

Real amplitude

Radial phase

Closed lines: soliton identification

Weak vs Strong Interactions of DOSs

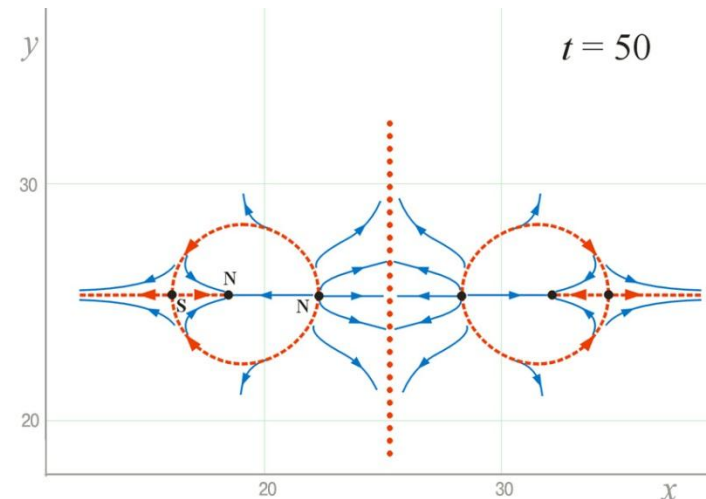
Weak interaction of DOSs: closed lines in the energy flow portraits are the same as for individual DOSs.

Strong interaction of DOSs: disappearance of some closed lines of individual DOSs and/or birth of new lines common for a number of DOSs.

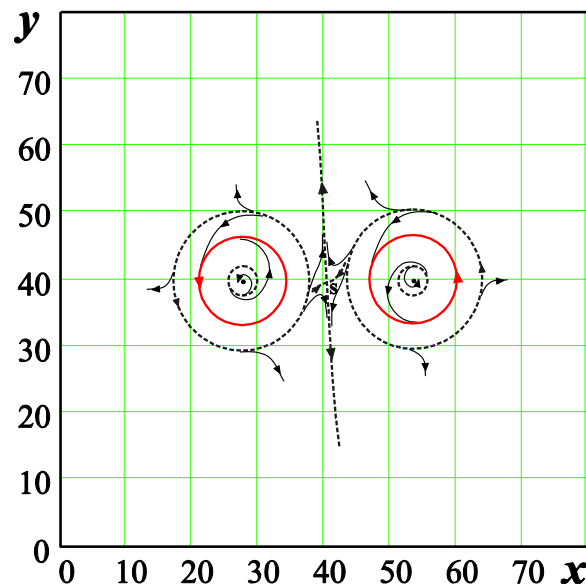
Energy flow portraits for laser DOSs

N – node, S – saddle

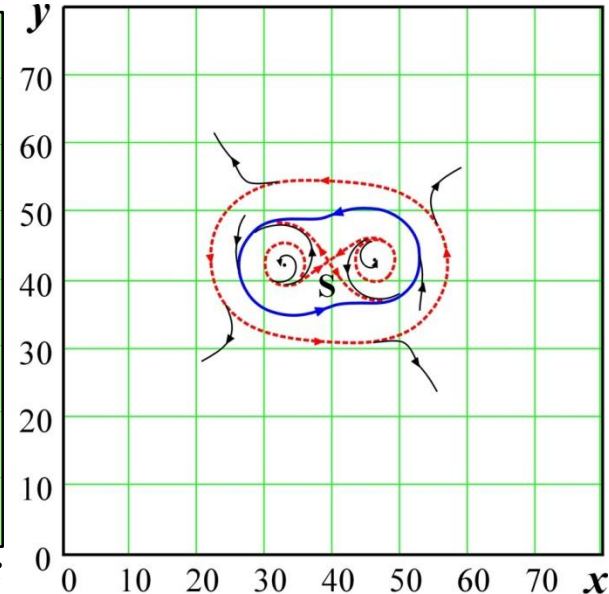
$t = 50$



Weak interaction of two fundamental DOSs



Weak interaction of two vortex DOSs

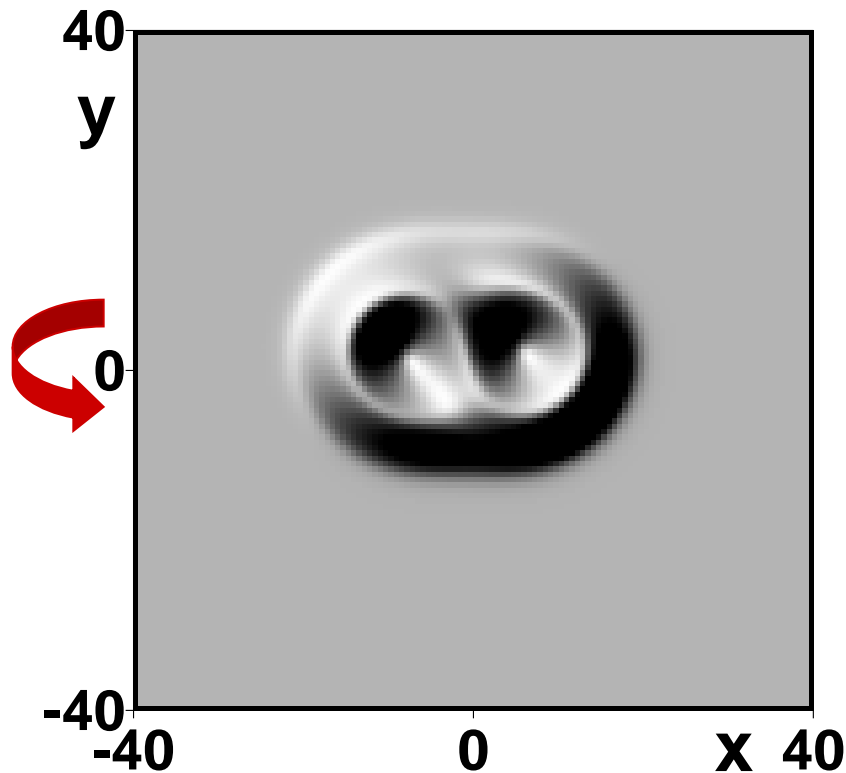


Strong interaction of two vortex DOSs

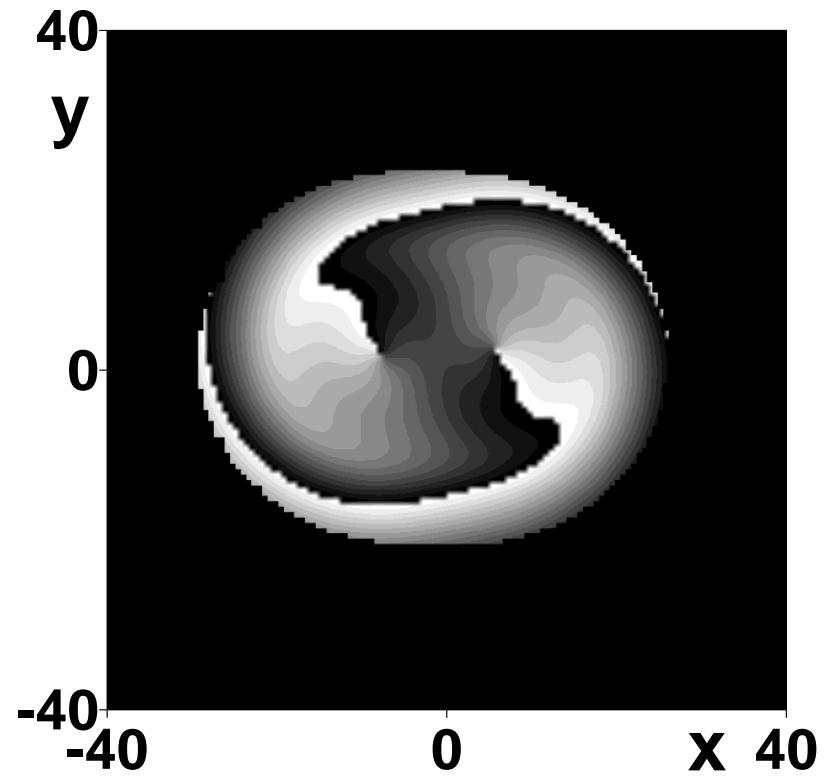
Rotating Pair of *STRONGLY* Coupled Vortex Solitons, $m = 1$

Central symmetry. Rotation, no revolution

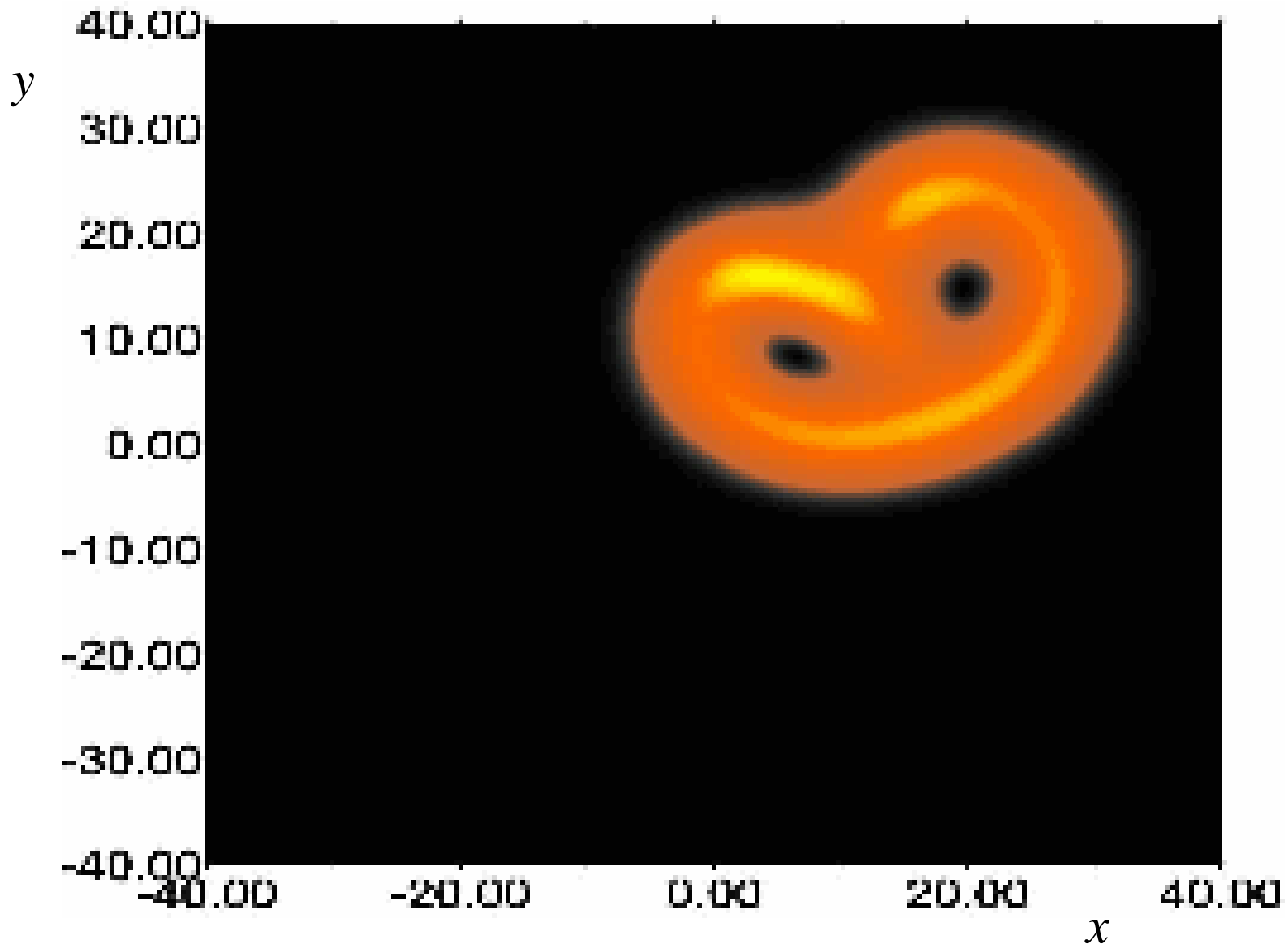
Intensity



Phase

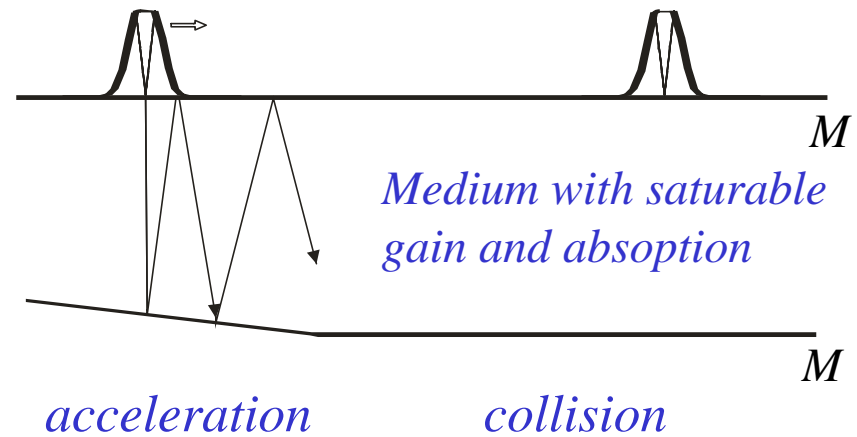


Two-Vortex Laser Soliton*



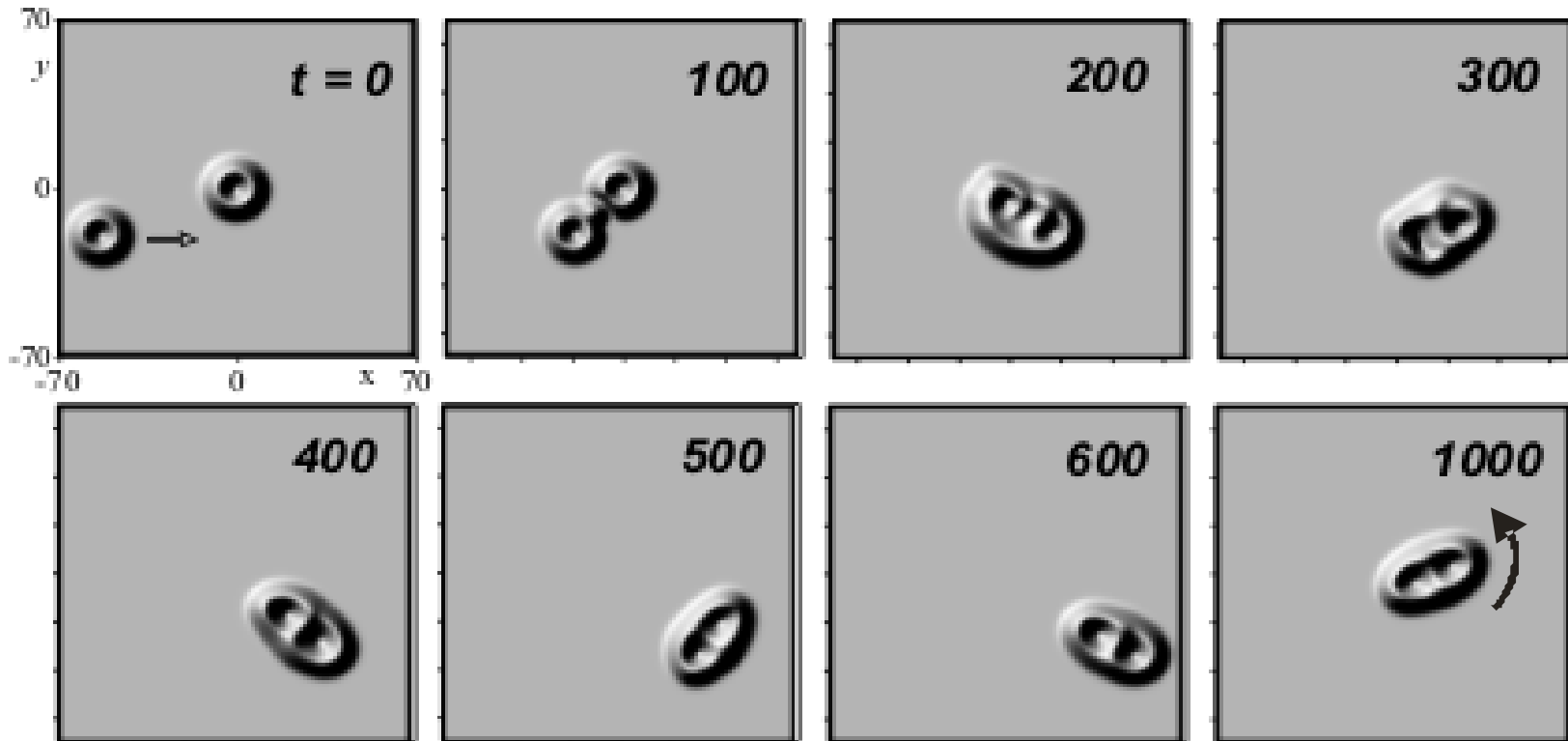
Formation of Strong Coupling

“Soliton collider”



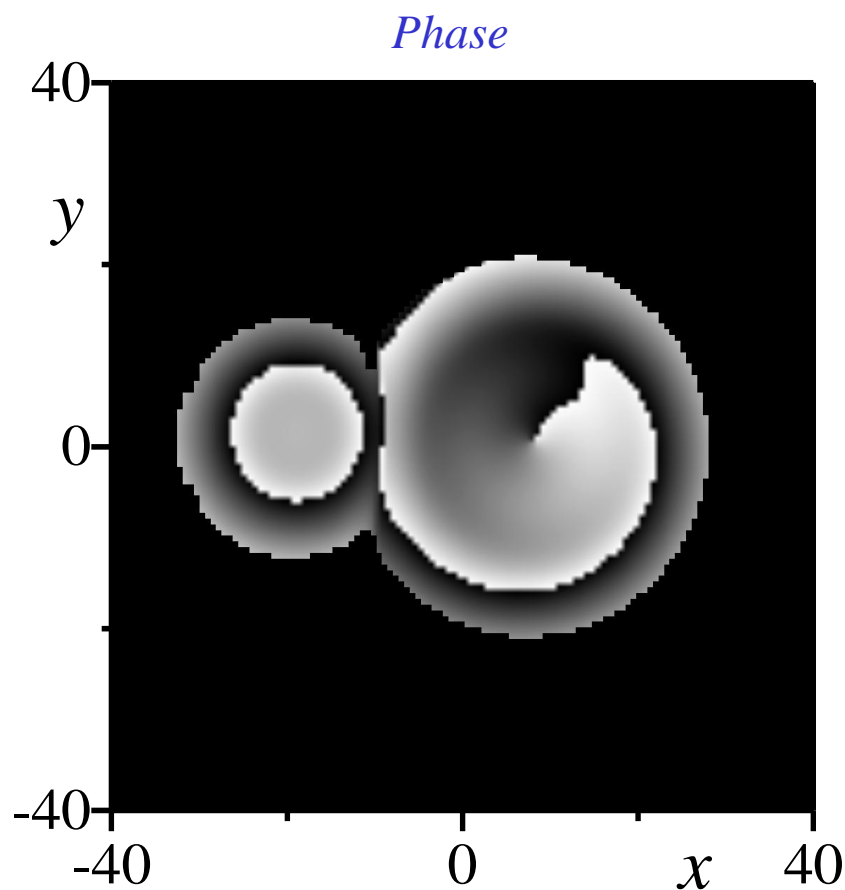
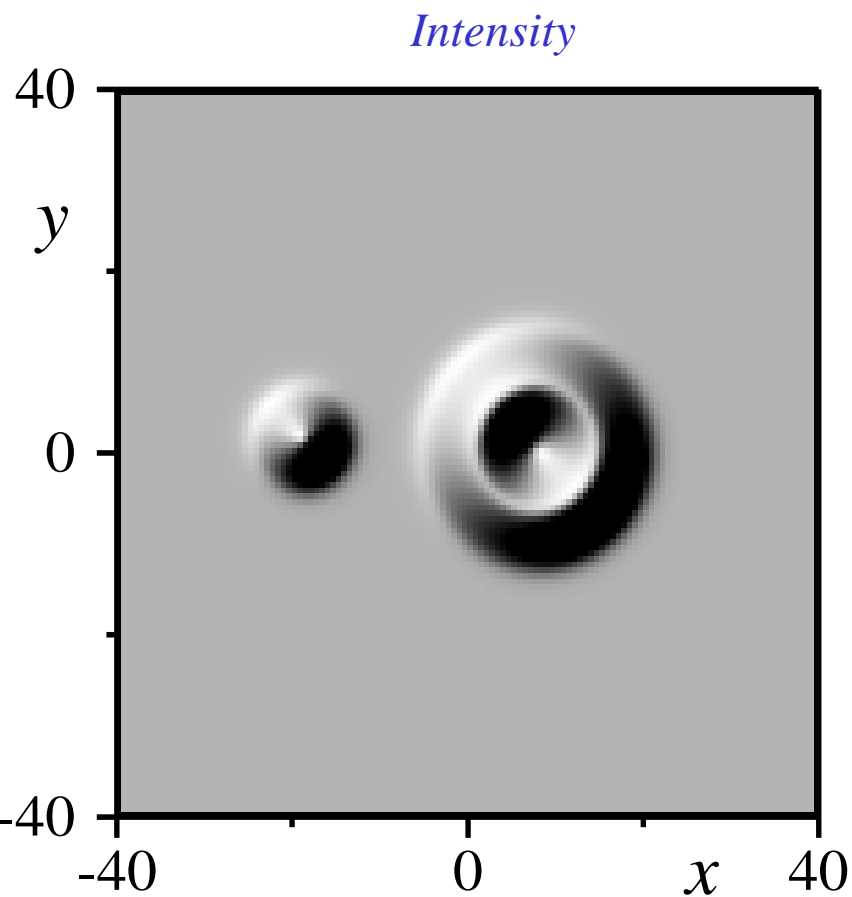
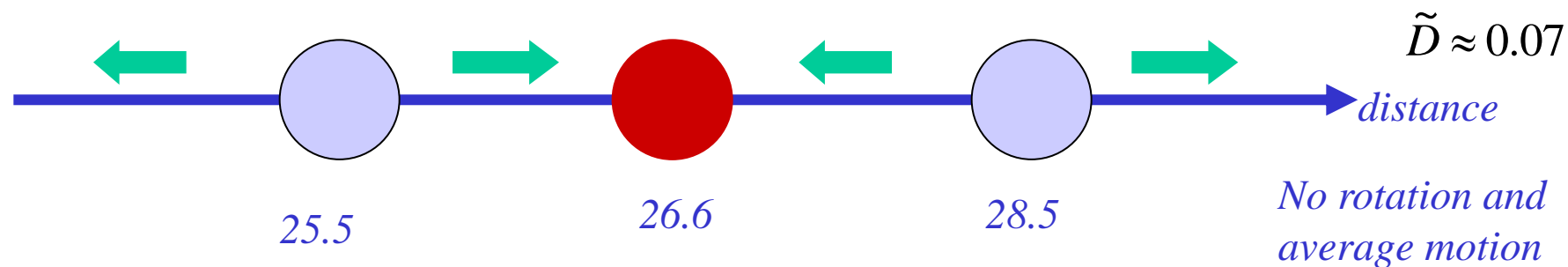
M - mirror

Intensity profiles



Incoherent coupling: Two laser DOSs with different topological charges and frequencies

$$m_1 = 0, m_2 = 1$$



Symmetry Elements

Simultaneous symmetry of instantaneous transverse distributions of intensity and energy flows

1. Axes of symmetry
2. Symmetry to rotation to angle $2\pi/n$, $n = 2, 3, \dots$

We will consider the cases when the symmetry is stable

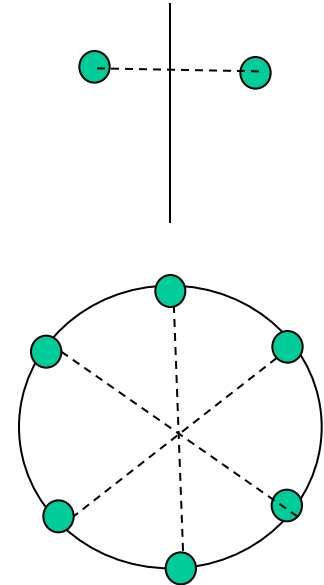
Below: Four variants of symmetry and motion of stationary “solid-like” soliton complexes [Rosanov 2007]

(Transverse)
Coordinates
of DOS's centre

$$\mathbf{R}_c(t) = \frac{\int \mathbf{r}_\perp |E(\mathbf{r}_\perp) - E_b|^2 d\mathbf{r}_\perp}{\int |E(\mathbf{r}_\perp, t) - E_b|^2 d\mathbf{r}_\perp}$$


E_b – background field, $E_b = 0$ for laser (not-driven) schemes

Centre velocity: $\mathbf{V}_c(t) = \frac{d}{dt} \mathbf{R}_c$

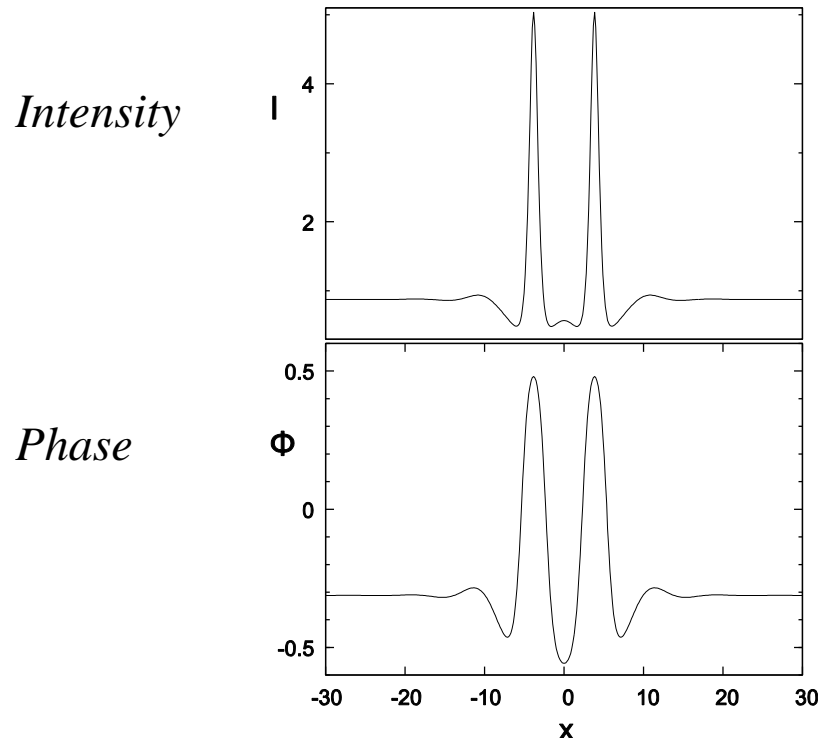


Eulerian Mechanics (solid-like structures).

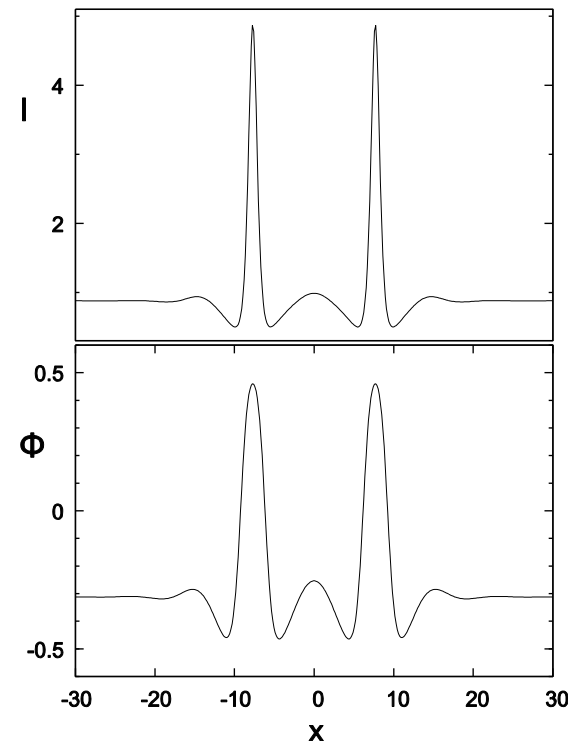
Driven 2D-interferometer (the same for lasers)

1. Two or more axes of symmetry  No (transverse) motion and rotation

Examples: single DOSs, pairs of DOSs with inter-soliton distances $d_{1,2}$; section $y = 0$



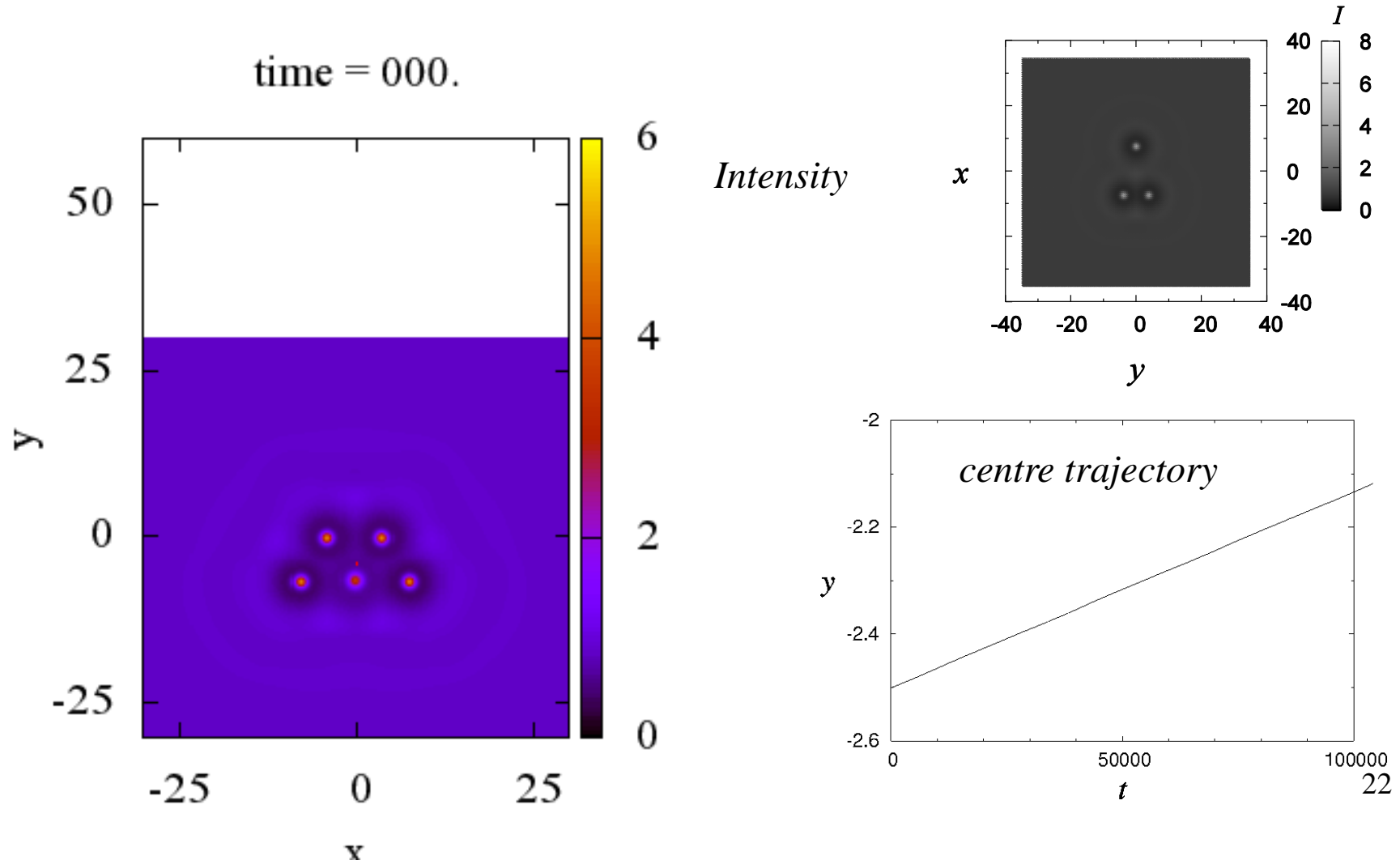
$d_1 = 7.66$



$d_2 = 15.43$

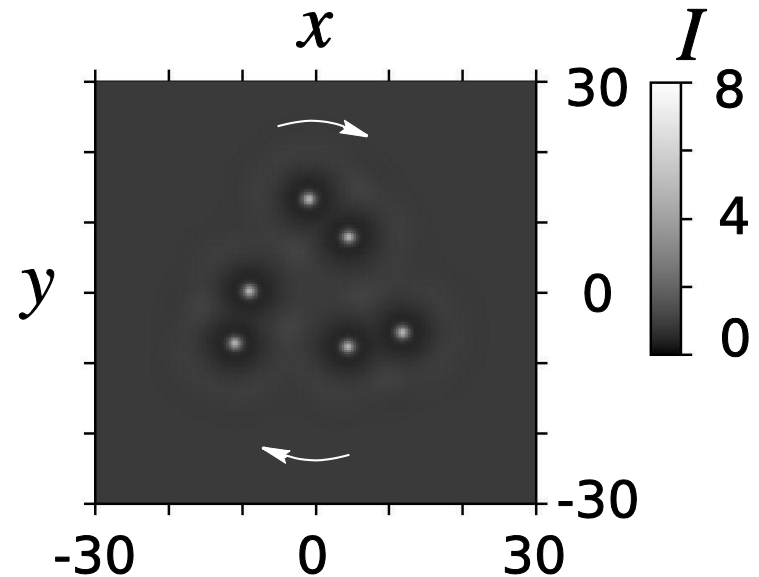
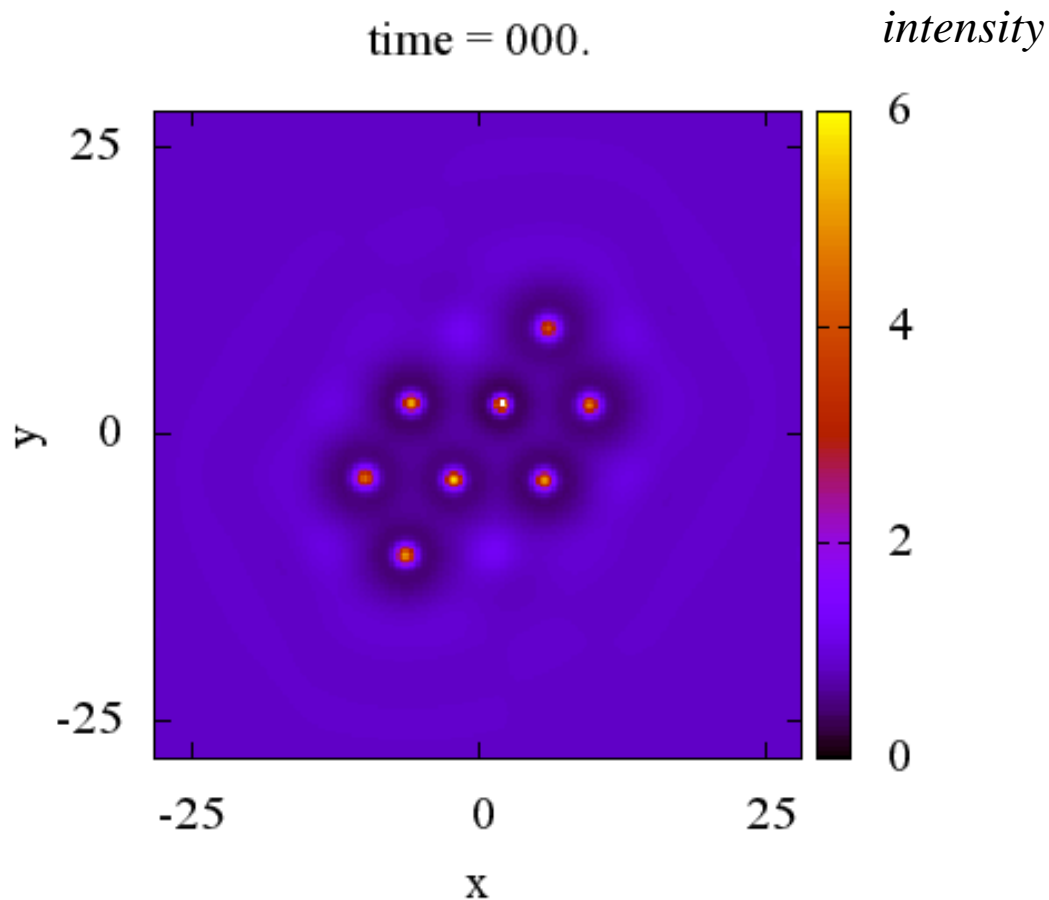
2. Only one axis of symmetry --- In-line motion with constant velocity without rotation

Driven nonlinear interferometer, weakly coupled DOSs

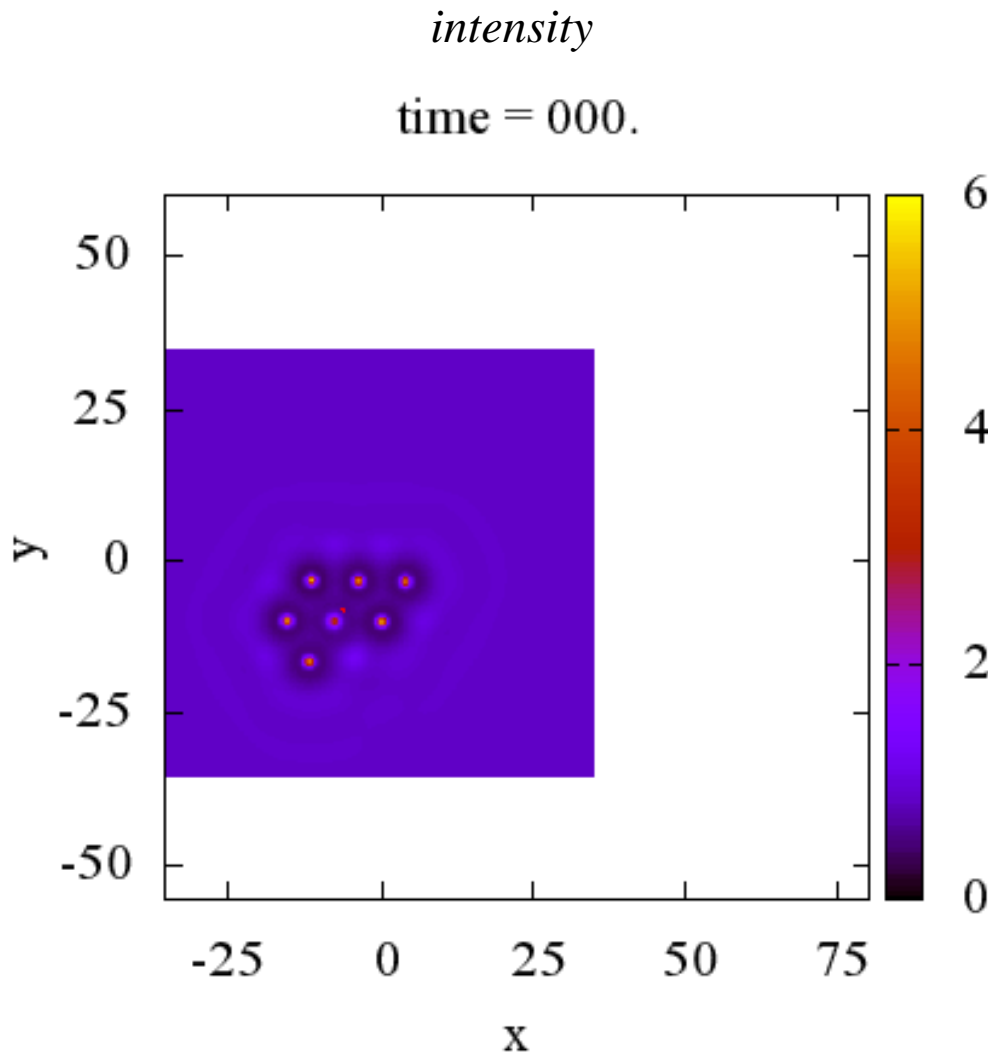


3. Symmetry to rotation --- Rotation with constant angular velocity, no centre motion

Driven nonlinear interferometer



4. No elements of symmetry - Moon-like motion

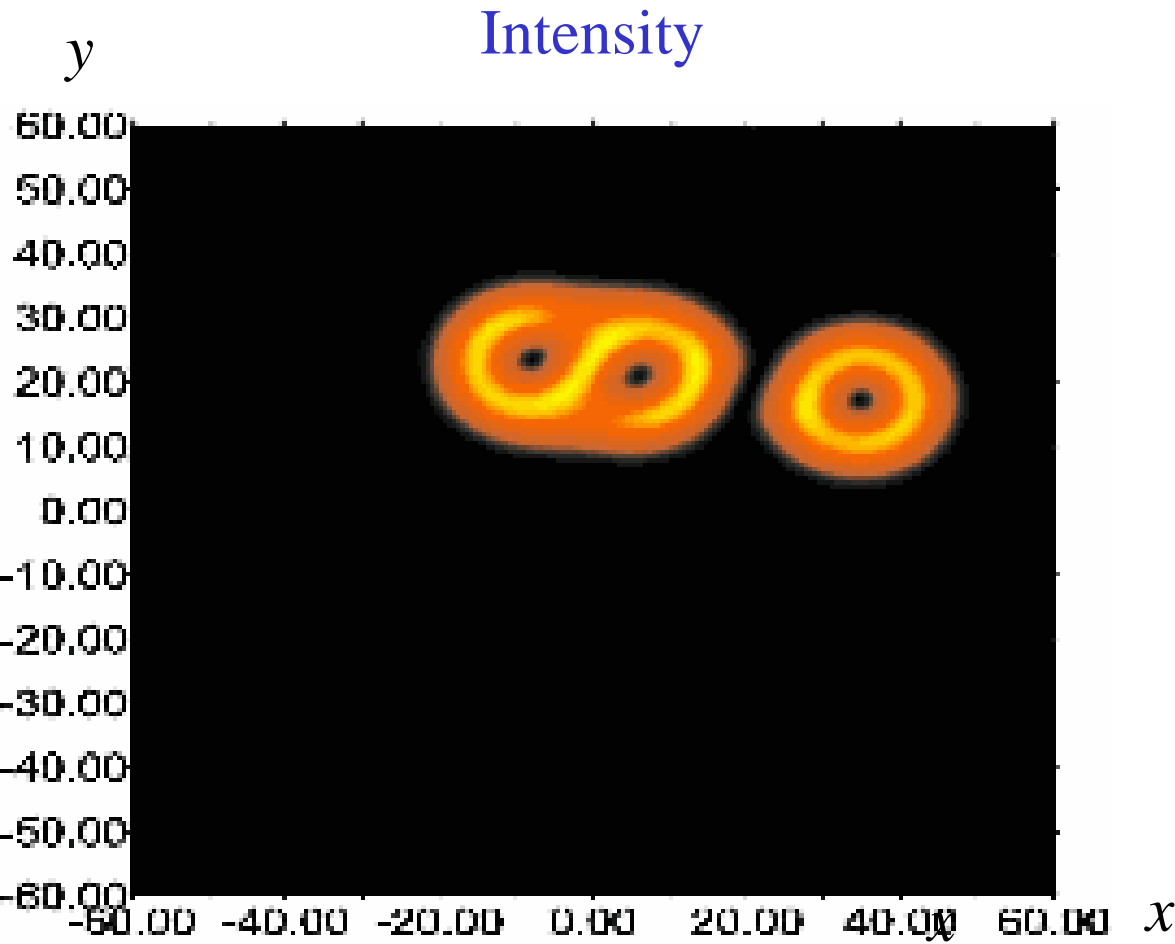


*Rotation of
asymmetric complex
with its simultaneous
rotation with the
same period*

More Complicated Motion of Laser Solitons: “Core” + “satellite”*

Topological charge $m = 1$

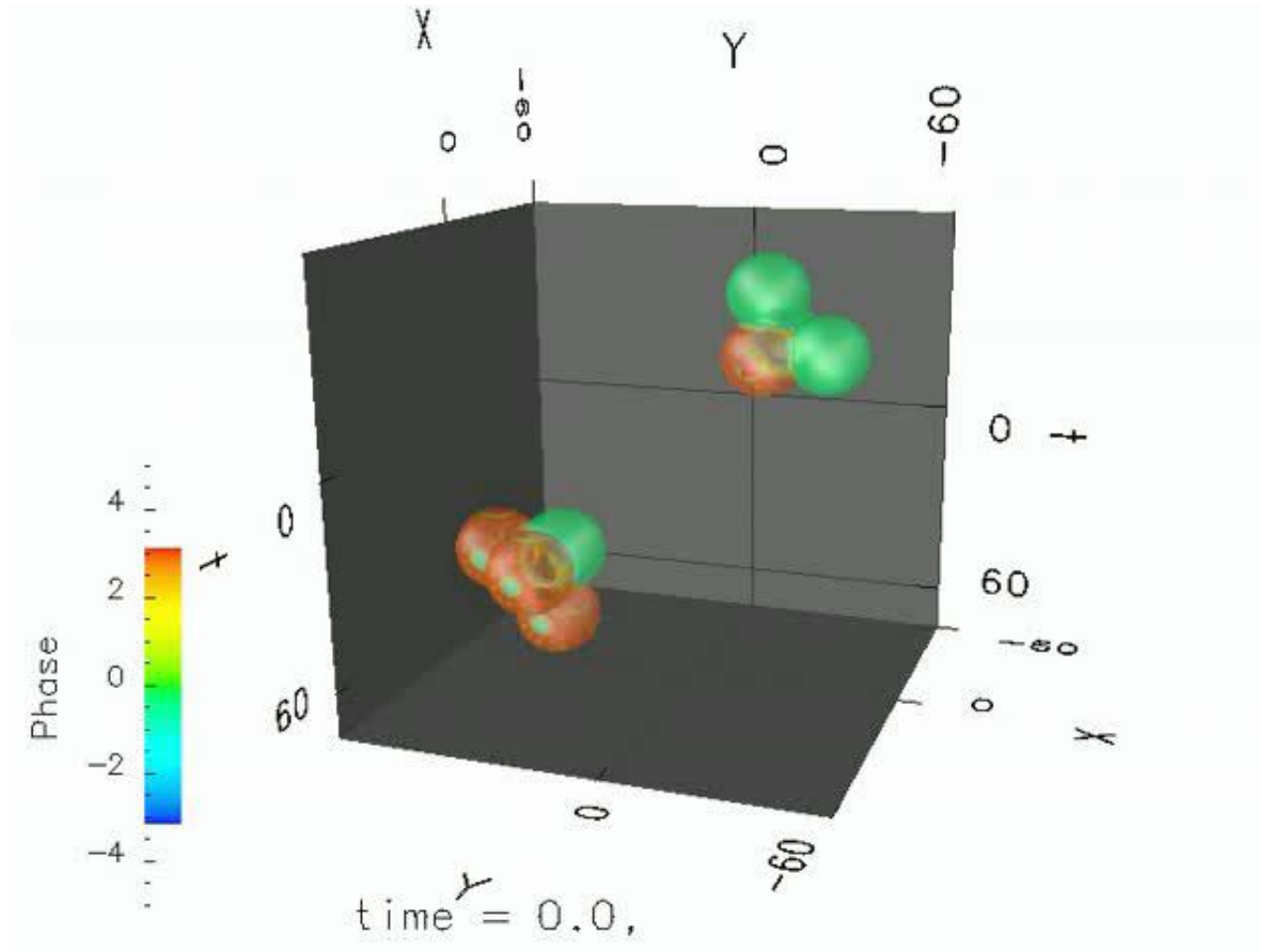
“Planetary system”



“Core” – a rotating pair of strongly coupled vortex solitons.

“Satellite” is weakly coupled with the core, their frequencies differ.

3D-Laser Soliton Complexes and Their Collision*

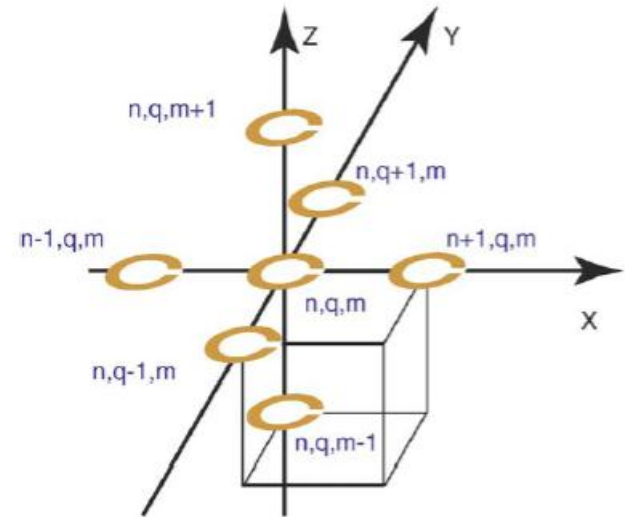
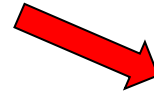


*Relative motion
in co-propagated
system of
coordinates*

Discrete Dissipative Solitons (*microwaves*)

A cubic lattice of weakly coupled nonlinear splitting resonators driven by coherent holding radiation

Holding radiation



$$\begin{aligned}
 i \frac{d}{d\tau} \Psi_{n,q,m} - (2\Omega - i\gamma + \alpha |\Psi_{n,q,m}|^2) \Psi_{n,q,m} - \Sigma = \\
 2\kappa (\Psi_{n,q,m+1} + \Psi_{n,q,m-1} - 2\Psi_{n,q,m}) - \\
 \kappa (\Psi_{n+1,q,m} + \Psi_{n-1,q,m} + \Psi_{n,q+1,m} + \Psi_{n,q-1,m} - 4\Psi_{n,q,m})
 \end{aligned}$$

Ψ – magnetization, Ω – frequency detuning, γ – losses,
 Σ – holding radiation amplitude, κ – coupling coefficient
Analytical perturbation approach for weak coupling

Discrete Dissipative Solitons, 1D and 2D



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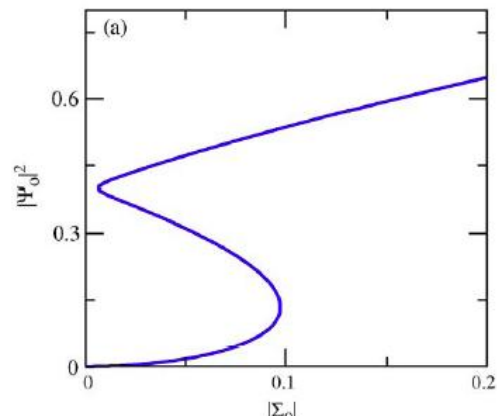
Photonics and Nanostructures – Fundamentals and Applications 4 (2006) 69–74

www.elsevier.com/locate/photronics

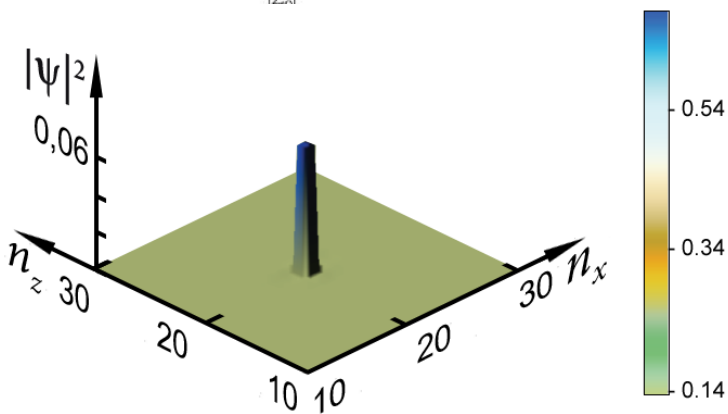
PHOTONICS AND
NANOSTRUCTURES
Fundamentals and Applications

Nonlinear magnetoinductive waves and domain walls in composite metamaterials

Ilya V. Shadrivov^{a,*}, Alexander A. Zharov^{a,b}, Nina A. Zharova^{a,c},
Yuri S. Kivshar^a



Bistability

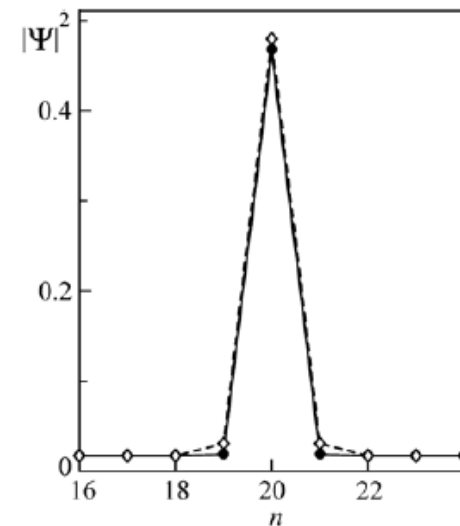


2D-discrete DOSs

ISSN 0021-3640/JETP Letters, 2011, Vol. 93, No. 12, pp. 743–746. © Pleiades Publishing, Inc., 2011.
Original Russian Text © N. N. Rosanov, N. V. Vysotina, A. N. Shatsev, I. V. Shadrivov, Yu. S. Kivshar, 2011, published in *Fizika i Zhurnal Eksperimental'noi i Teoreticheskoi Fiziki*, 2011, Vol. 93, No. 12, pp. 626–629.

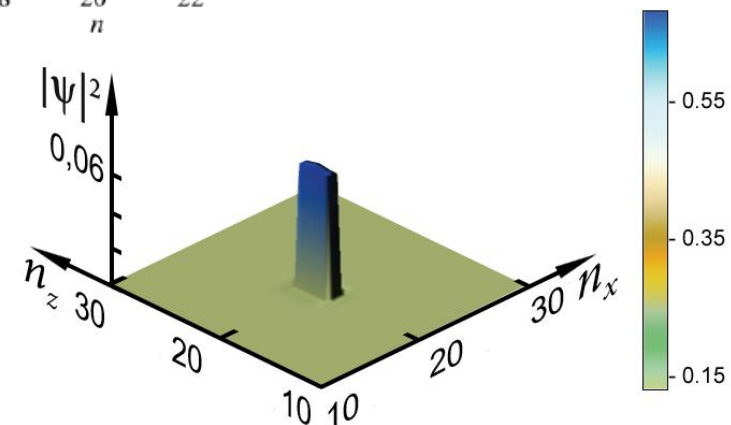
Hysteresis of Switching Waves and Dissipative Solitons in Nonlinear Magnetic Metamaterials

N. N. Rosanov^{a,b}, N. V. Vysotina^{a,b}, A. N. Shatsev^{a,b}, I. V. Shadrivov^{a,c}, and Yu. S. Kivshar^{a,c}



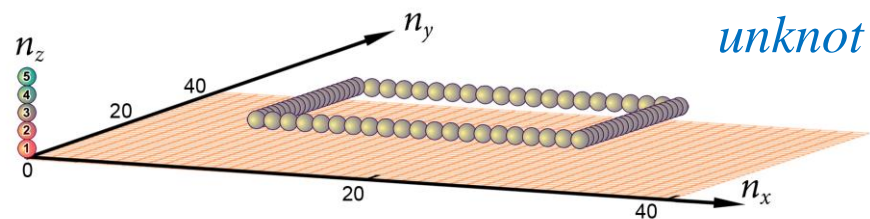
Perturbation theory with weak coupling

1D-discrete DOS

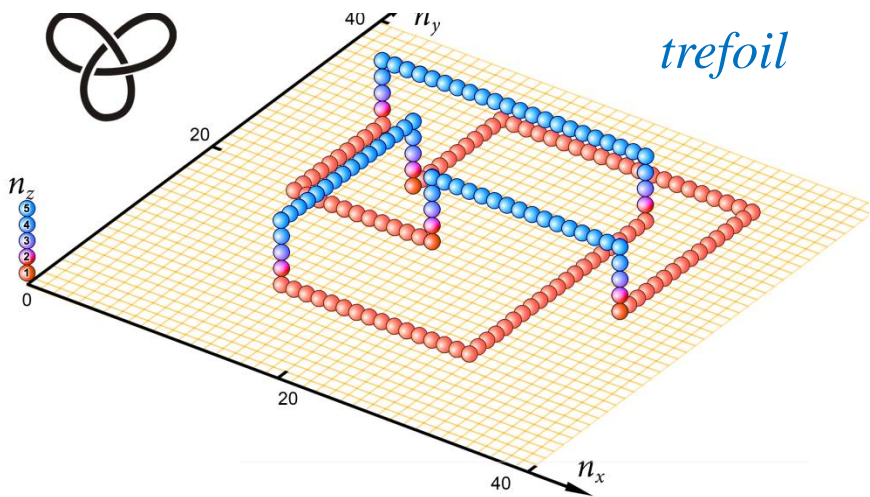
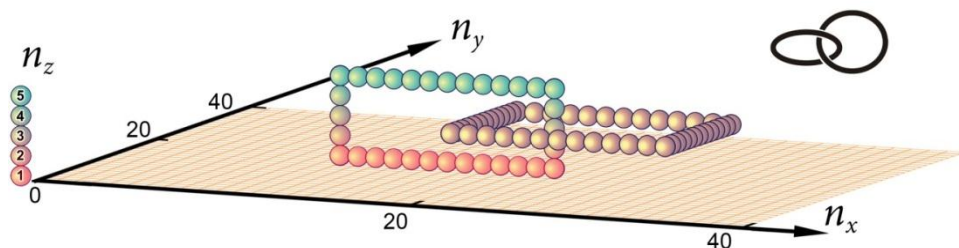


3D-Topological Discrete Solitons

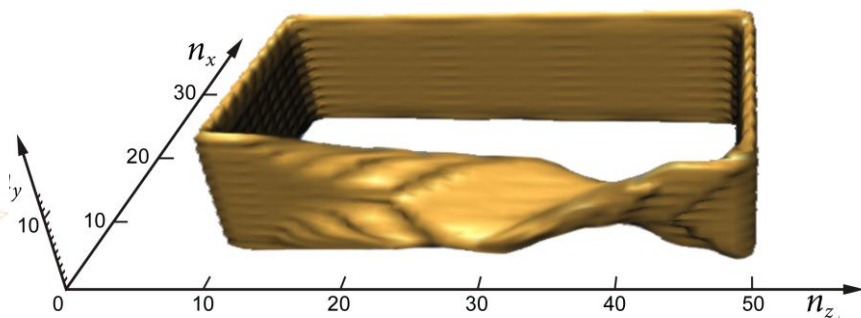
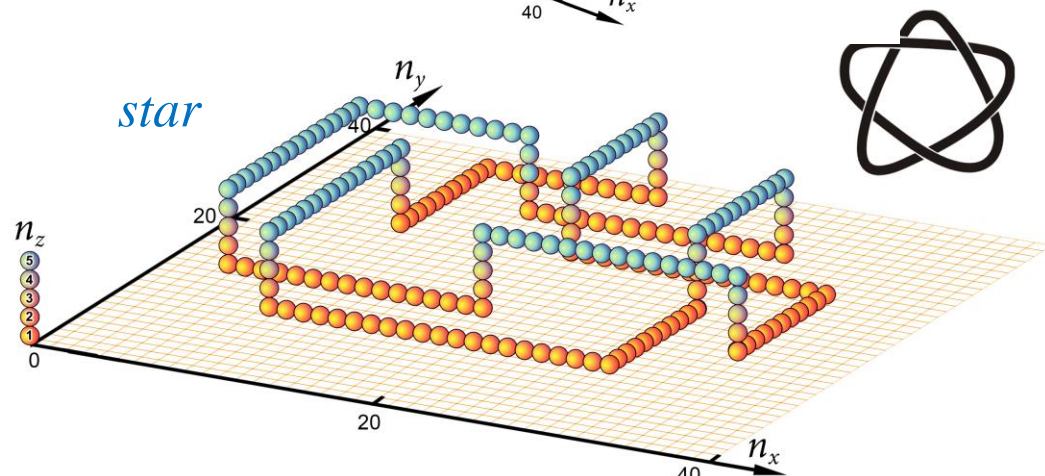
*Knots - closed lines without intersections –
location of elements excited to the upper state*



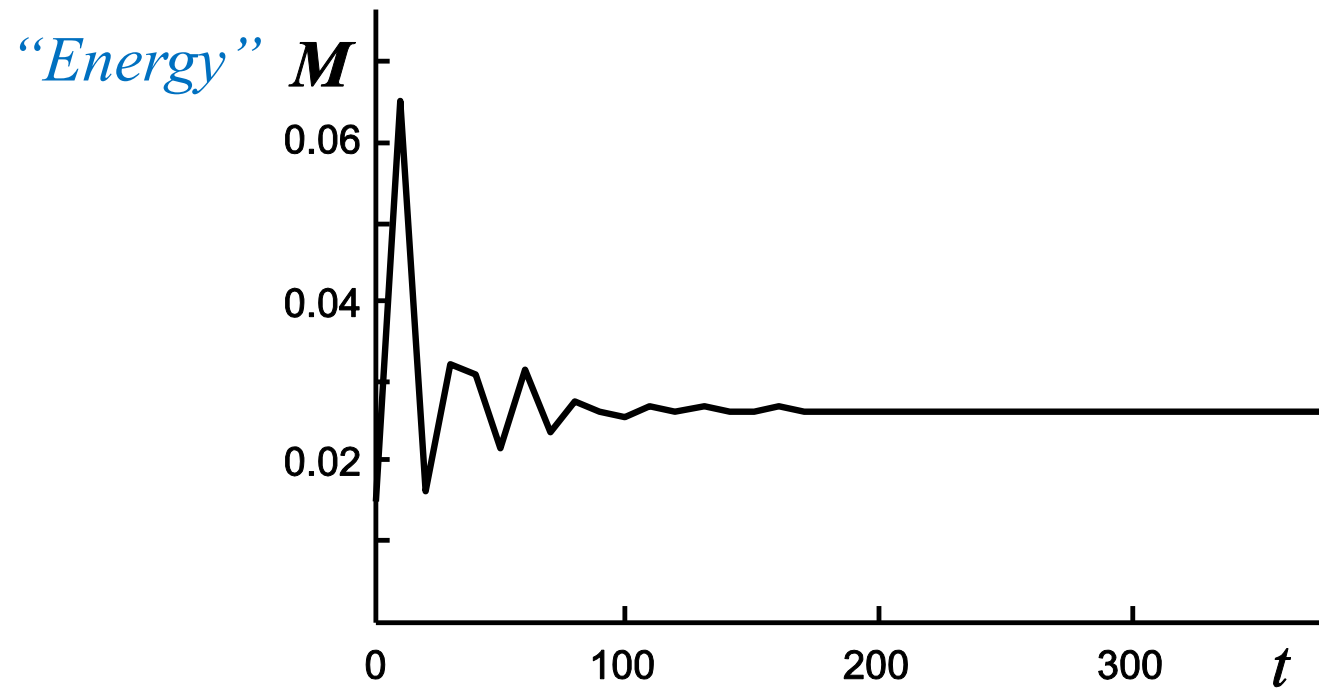
Hopf link



Möbius-like band



Stability

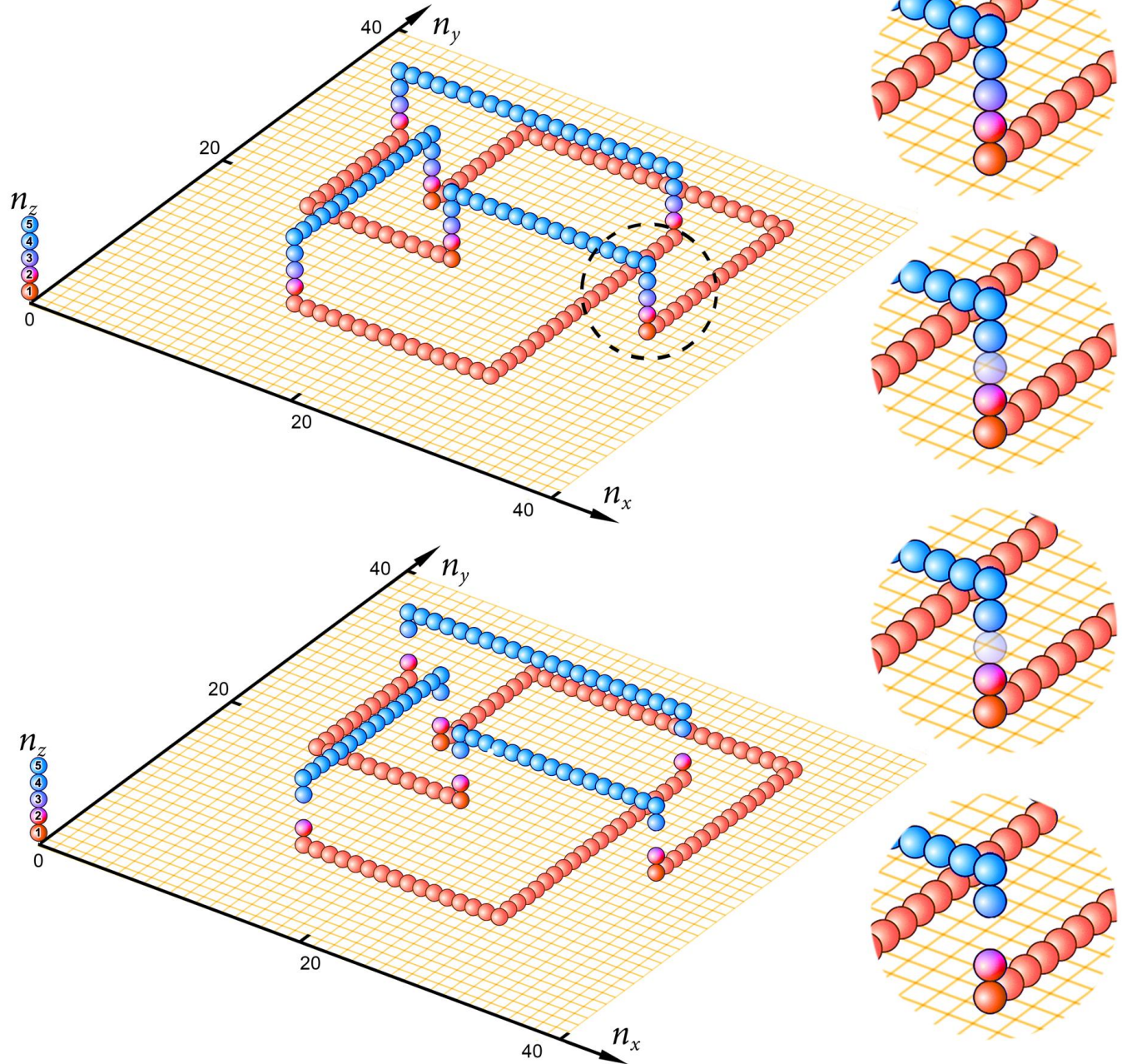


Transient period

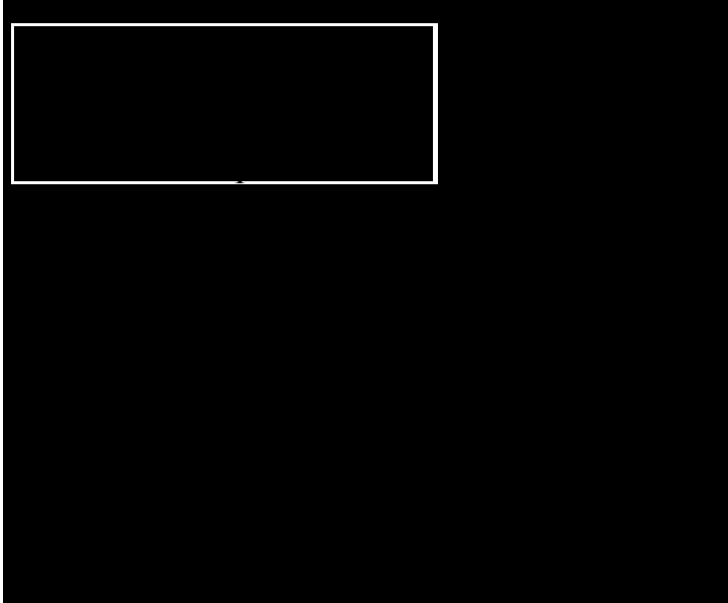
Decay of trefoil with increase of coupling coefficient

The scheme is highly anisotropic

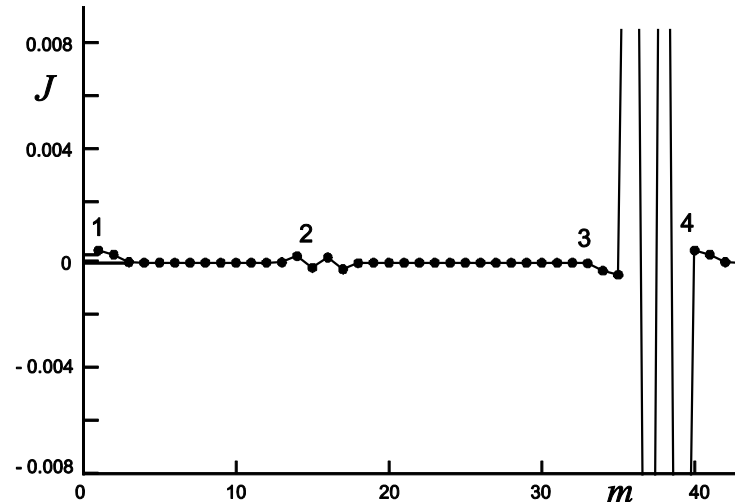
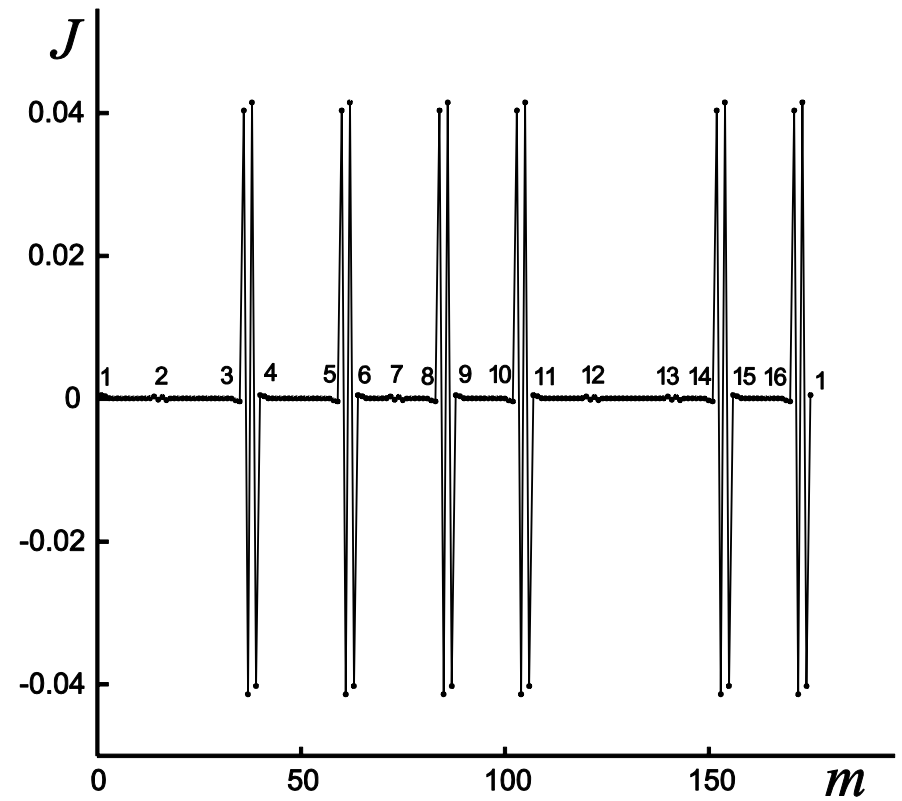
Line disintegration



Discrete Currents



$$J_m = 2 \operatorname{Im}(\Psi_m^* \Psi_{m+1})$$



Experiments

T. Elsass, K. Gauthron, G. Beaudoin, I. Sagnes, R. Kuszelewicz, S. Barbay. Fast manipulation of laser localized structures in a monolithic vertical cavity with saturable absorber. Appl. Phys. B 98, 327 (2010).

VCSEL, gain – two quantum wells

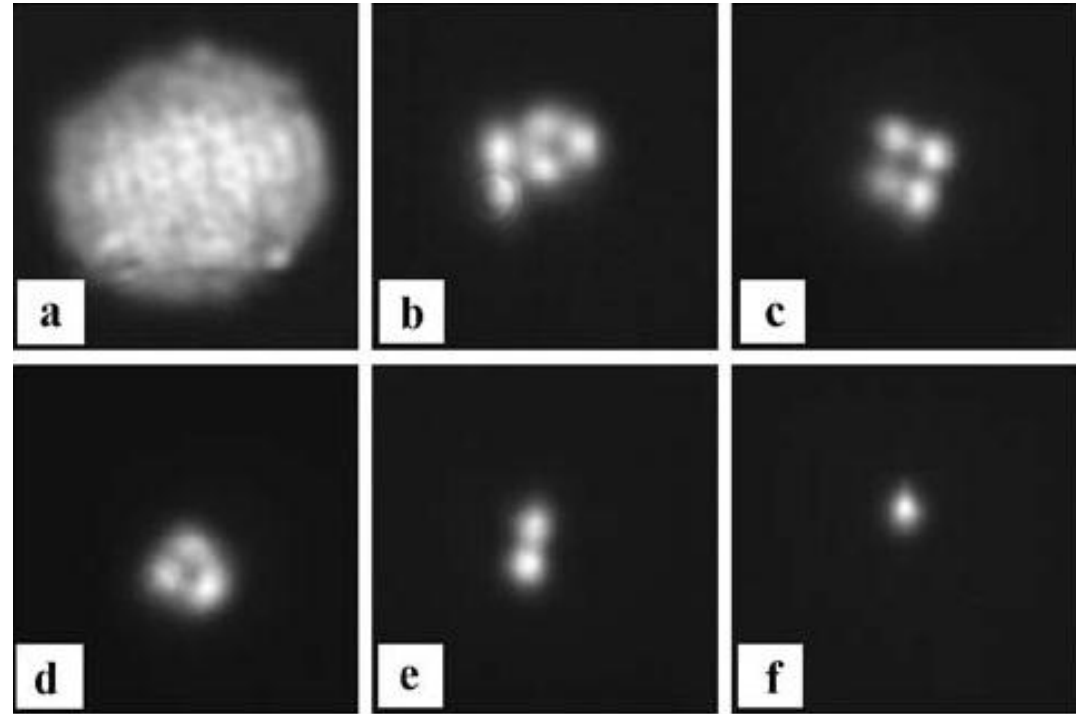
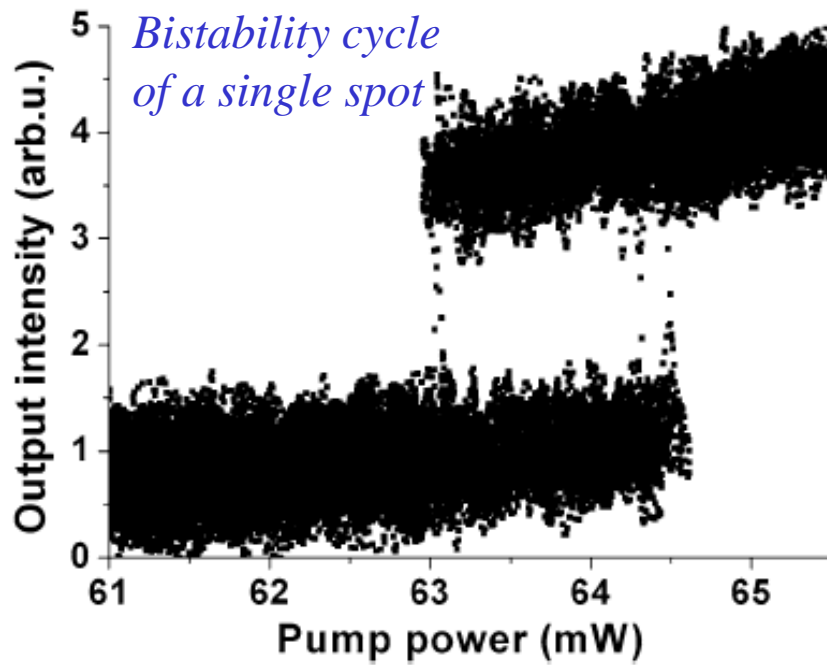
$\text{In}_{0.2}\text{Ga}_{0.8}\text{As}/\text{Al}_{0.05}\text{Ga}_{0.95}\text{As}$.

A saturable absorber: one quantum well

$\text{In}_{0.2}\text{Ga}_{0.8}\text{As}/\text{Al}_{0.22}\text{Ga}_{0.78}\text{As}$

Pump power = 20 – 130 mW

Writing process ~ 10 ns, erasure < 1 ns



*a – lasing over the whole aperture
b-e – complexes of DOSs
f – a single DOS*

Promising scheme: quantum dot laser with saturable absorber.

Extreme DOSs. Nanosolitons

DISSIPATIVE MOLECULAR SOLITONS IN LINEAR ORIENTED MOLECULAR J-AGGREGATES DRIVEN WITH LASER RADIATION

ISSN 0021-3640, JETP Letters, 2008, Vol. 87, No. 12, pp. 663–666. © Pleiades Publishing, Ltd., 2008.

Original Russian Text © Al.S. Kiselev, An.S. Kiselev, N.N. Rozanov, 2008, published in *Pis'ma v Zhurnal Éksperimental'noi i Teoreticheskoi Fiziki*, 2008, Vol. 87, No. 12, pp. 763–766.

Nanosized Discrete Dissipative Solitons in Resonantly Excited Molecular J-Aggregates

Al. S. Kiselev^{a,b}, An. S. Kiselev^{a,b}, and N. N. Rozanov^{a,b}

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^b *Vavilov State Optical Institute, St. Petersburg, 199034 Russia*

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Received May 5, 2008

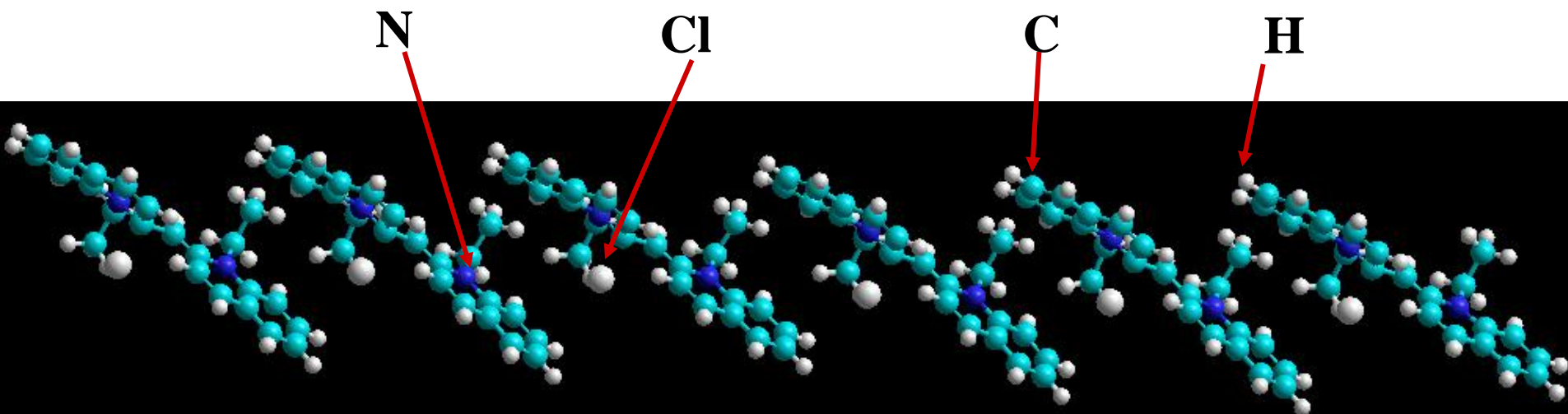
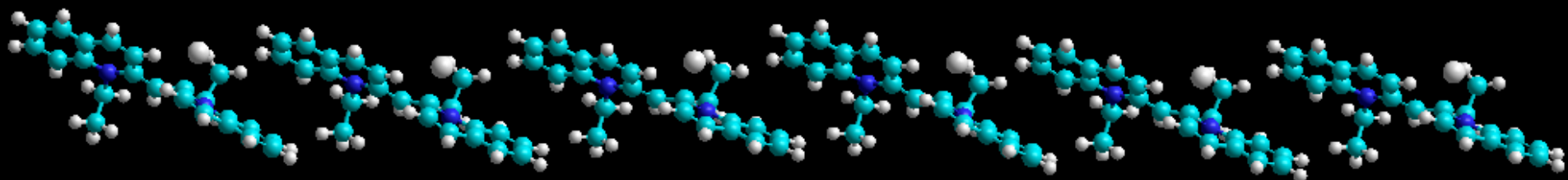
The resonant excitation of a linear oriented molecular chain simulated by a system of two-level schemes interacting through radiation has been analyzed. The regimes of modulation instability, switching waves, and dissipative solitons whose sizes for J-aggregates can reach about 1 nm have been revealed.

PACS numbers: 36.40.Vz, 42.75.Pc

Linear oligomers of molecules pseudoisocyanine-chloride $\text{PIC}_n:\text{Cl}_n$

$\text{PIC} = \textit{Pseudoisocyanine}$

Stable configuration:

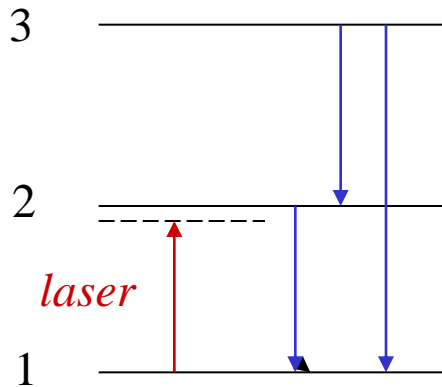


Quantum-chemistry code GAMES
[V.G. Maslov]

Metastable configuration

Simplified Model

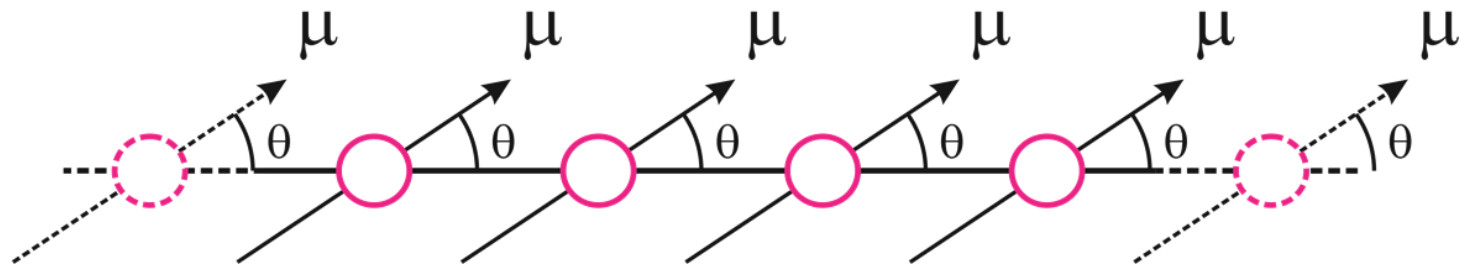
*One-particle density matrix approach; three-level scheme of electronic transitions;
bistability [Malyshev et al. 1998]*



*~ quantum-mechanical Neuman (Bloch) equations
for density matrix*

Interaction of molecules via radiation

$$E_n = E_{laser} + \sum_{m \neq n} E_{mn}$$



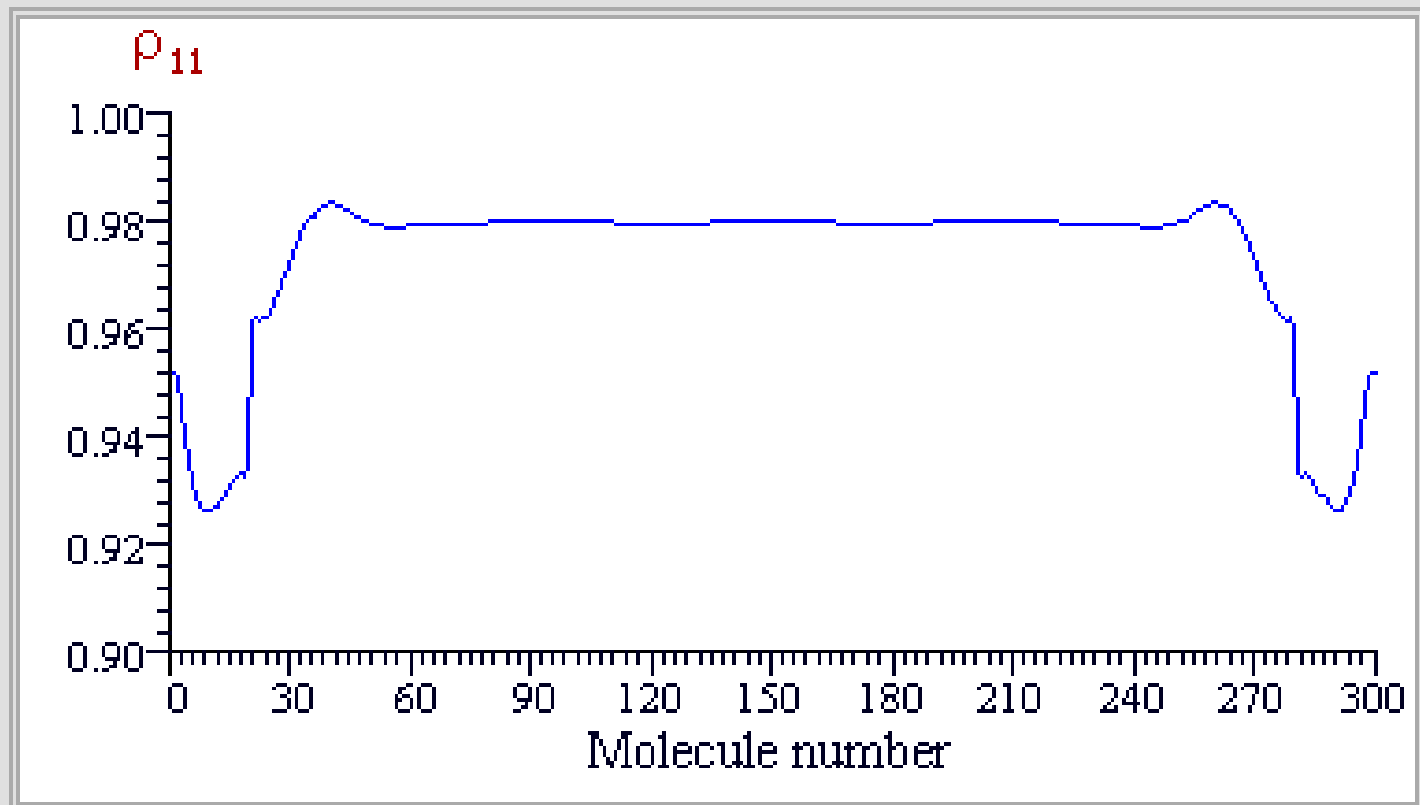
holding cw-laser radiation



Nanosoliton Formation*

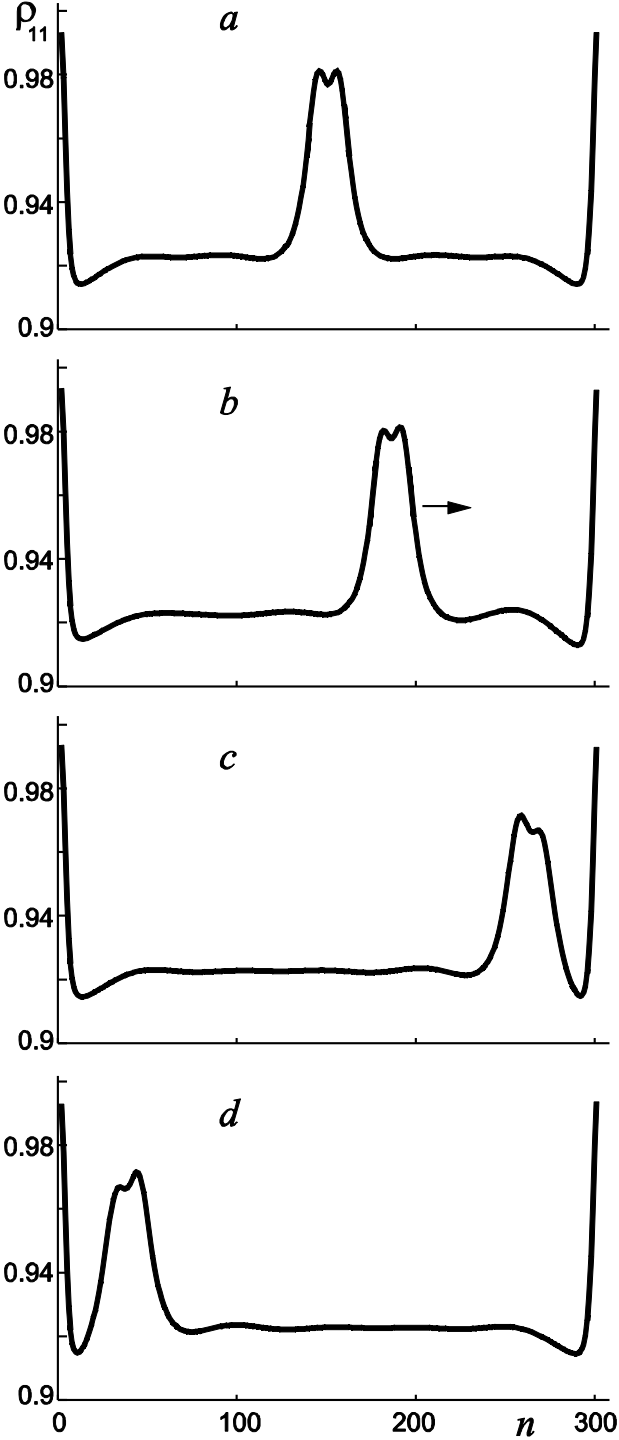
T = 1.000

J-AGGREGATE



Controlling Soliton Position

Oblique incidence of laser radiation



(a) angle of incidence $\varphi = 0$,
motionless soliton;

(b) $\varphi = 0.1$, soliton moves to the right;

(c) $\varphi = 0.1$, soliton stops
near the chain right edge;

(d) $\varphi = -0.1$, soliton moves to the left
and stops near the chain left edge.

**Soliton can be moved through
the whole molecular chain
in both directions**

Towards Attosolitons:

Few-cycle dissipative optical solitons

Idea: Self-induced transparency, effect of dissipative factors

Scheme: single-mode fiber doped with active (with pump) and passive (without pump) centres (atoms, ions, quantum dots).

Without active centres and dissipative factors, conservative solitons of self-induced transparency form a family with continuously varying parameter (maximum intensity and corresponding width).

After introduction of weak linear small-signal gain, there is transition in the family to solitons with higher intensity and smaller width (collapse).

When any mechanism arrests the collapse, extremely short localized pulses with extremely high peak intensity should occur.

N. V. Vyssotina, N. N. Rosanov, V. E. Semenov, S. V. Fedorov, S. Wabnitz, 2006

Model

Radiation propagation: 1D (plane waves), fixed polarization

Full wave equation:

$$\frac{\partial^2 E}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} (E + 4\pi P) = 0$$

E – electric field

t – time,

z – propagation coordinate

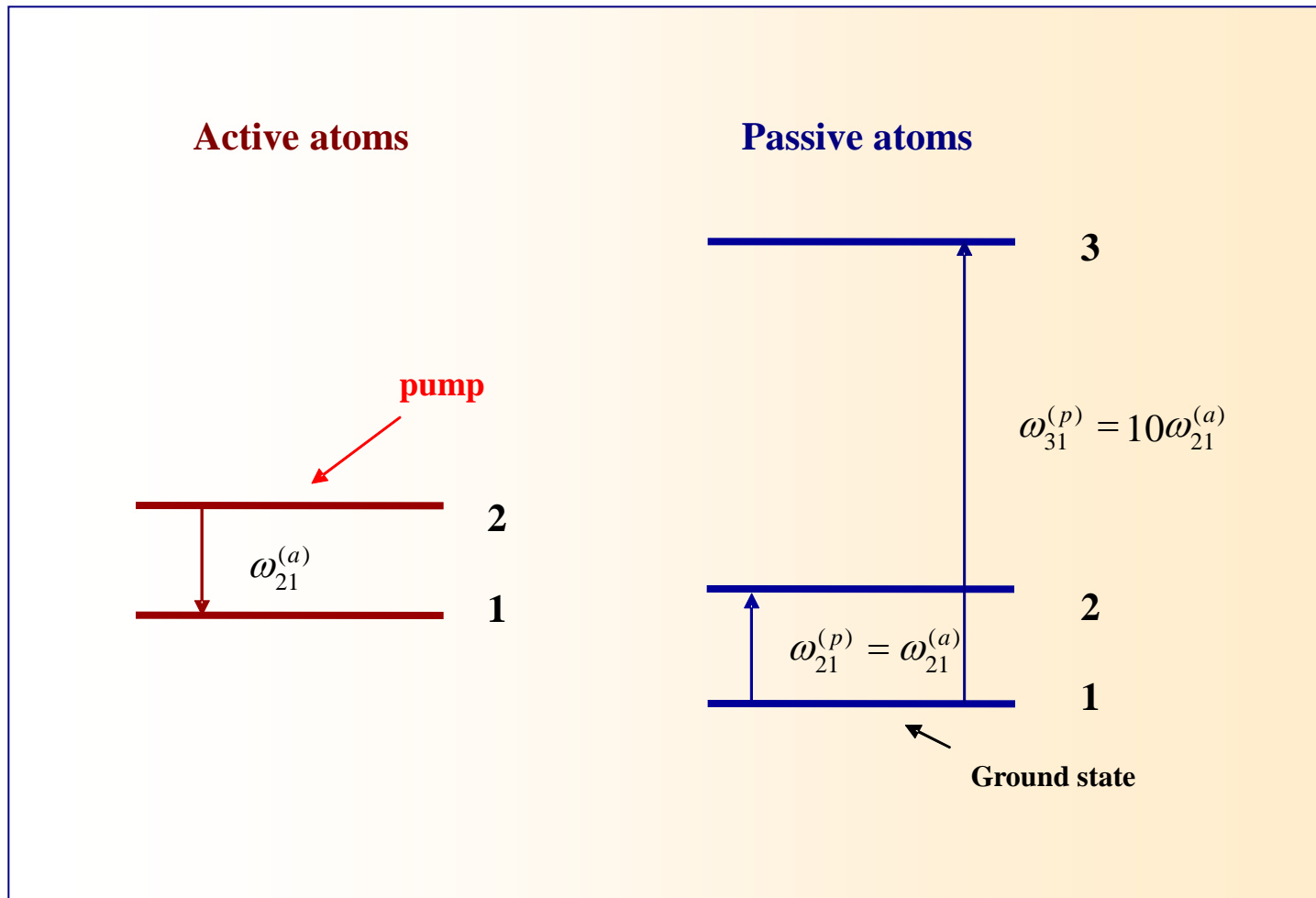
c – light velocity in vacuum

P – medium polarization

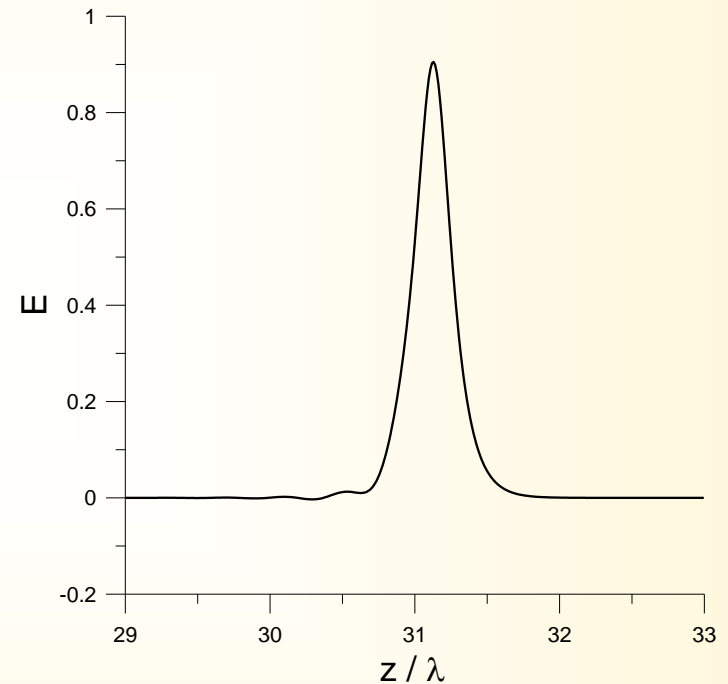
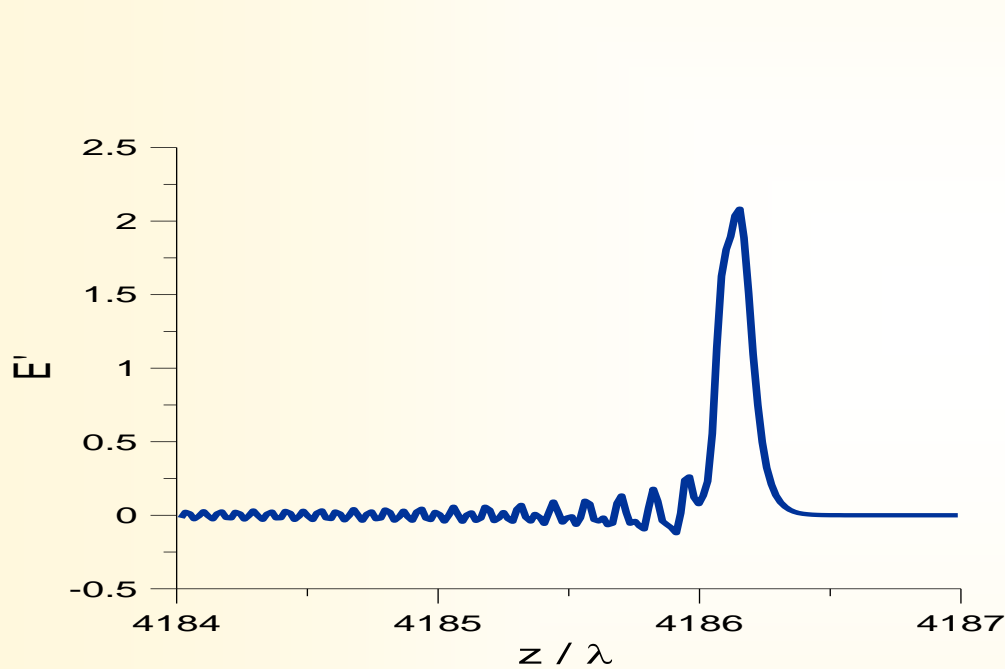
No approximations of slowly varying
amplitude and unidirectional propagation

Constitutive equations – quantum mechanical
Neuman (Bloch) equations for density matrix

Variant: Three-level passive atoms + two-level active atoms

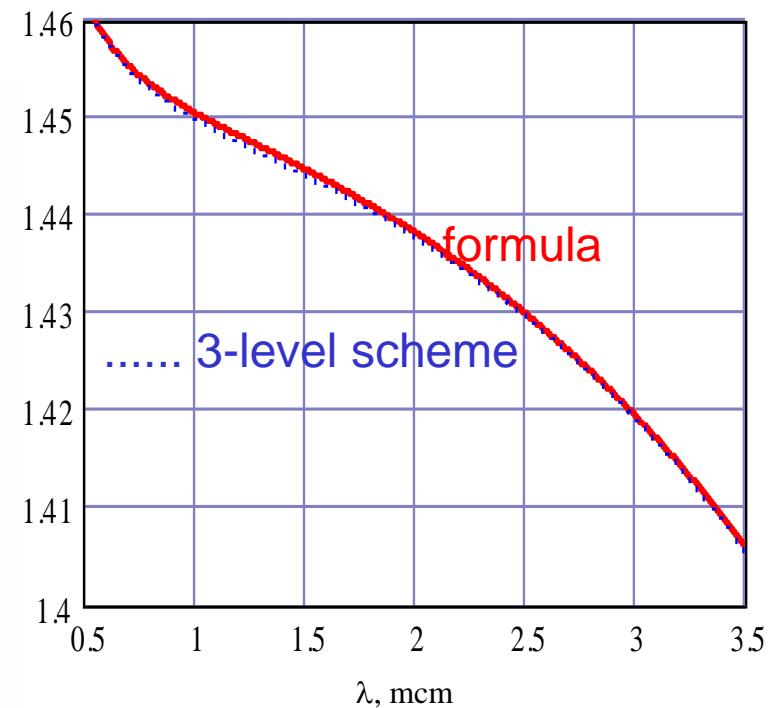
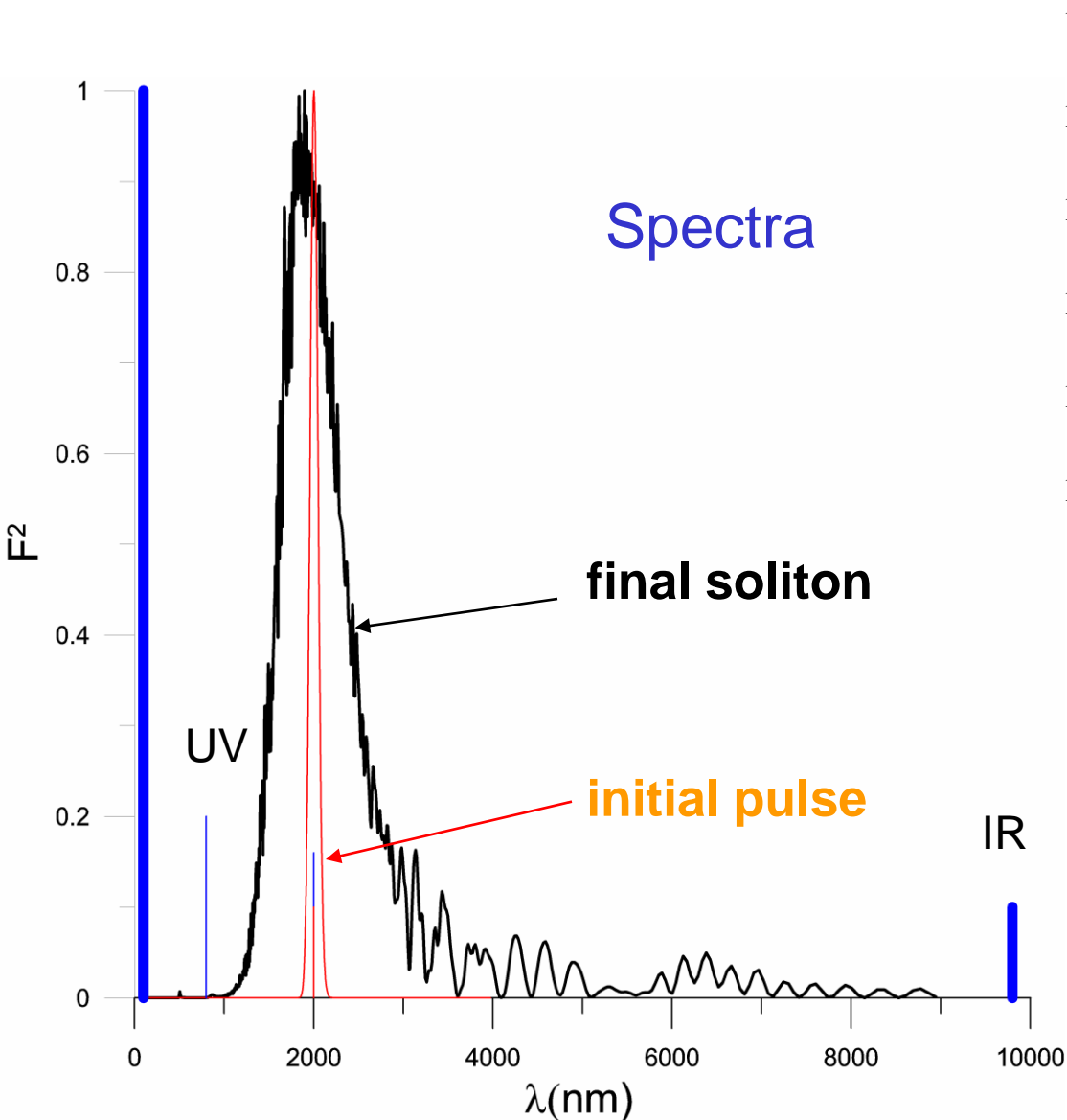


Examples of stable video-solitons (full field, not envelope)



(for different parameters of the 3rd level)

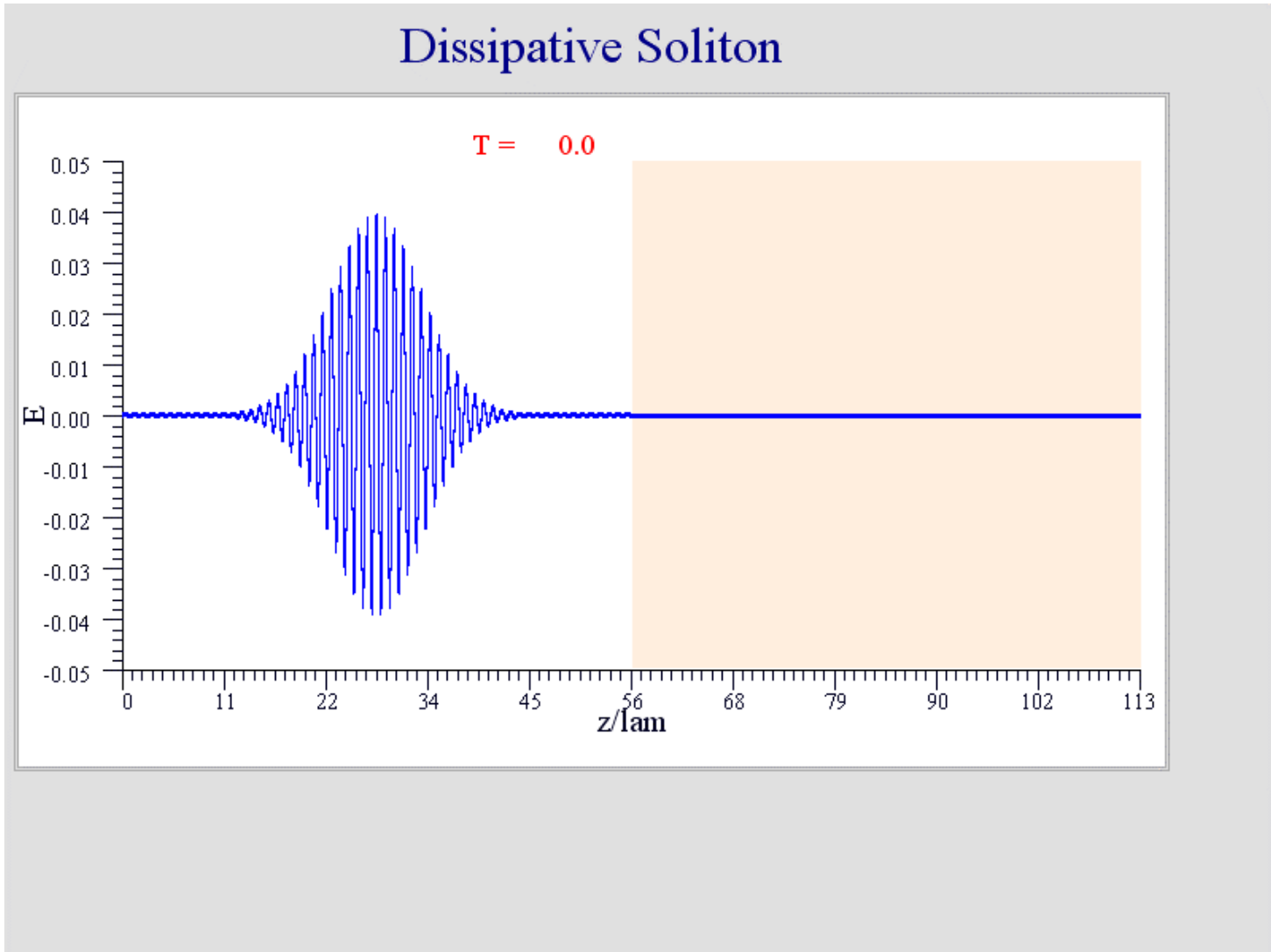
Taking into account matrix absorption bands



Refractive index in the
transparency band

*Coherent supercontinuum.
Promising are schemes with
quantum dots*

Formation of Few-cycle DOS from an Initial Femtosecond Pulse*



Conclusions

Standard, or envelope Dissipative Optical Solitons (DOSs):

- Quasi-quantum features – *discrete type of the spectrum of main parameters* – provide *higher stability* of DOSs in the presence of noise. Therefore they are promising for information applications with heightened requirements to the precision and operation reliability.
- *Large diversity* of DOSs' types. They are not “elementary”, but have *internal structure* determined by the *topology of energy flows*. This topology allows one to distinguish between *weak and strong interactions* of DOSs.
- DOS mobility and mechanics. *Symmetry and Eulerian mechanics* of solid-like soliton complexes. Their *curvilinear motion*.
- Not presented here topics: DOSs' quantum features (Brownian motion and squeezed states); Effect of relaxation; ...

Tendency to extreme (not envelope) solitons.

Nanosized dissipative molecular solitons.

Few cycle solitons using dissipative (gain/absorption) factors. Nonlinear widening of gain spectrum.

Not achieved yet: Combination of nanosizes and attoduration → light dots

Great scope for fundamental research and applications.

Conclusion in Short:

Brevius, Angustius, Fortius!

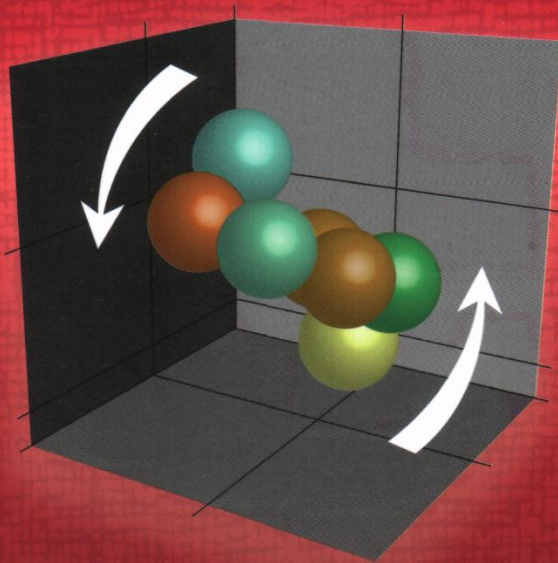
(Shorter – in time and space, Stronger!)

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Н. Н. РОЗАНОВ

ДИССИПАТИВНЫЕ ОПТИЧЕСКИЕ СОЛИТОНЫ

ОТ МИКРО-
К НАНО-
И АТТО-



N.N. Rosanov.
Dissipative
Optical Solitons.
From Micro- to
Nano- and Atto-
Moscow,
Fizmatlit, 2011
(in Russian)

15th International Conference
«Laser Optics 2012»

St. Petersburg, Russia, June 25-29, 2012

High-Field and High-Energy Lasers

Lightwave Communication Technologies

Biophotonics and Nanophotonics

Nonlinear Photonics and Metamaterials

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English will be the official language of the Conference

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